



Overview of DIS results, global fittings and DY predictions

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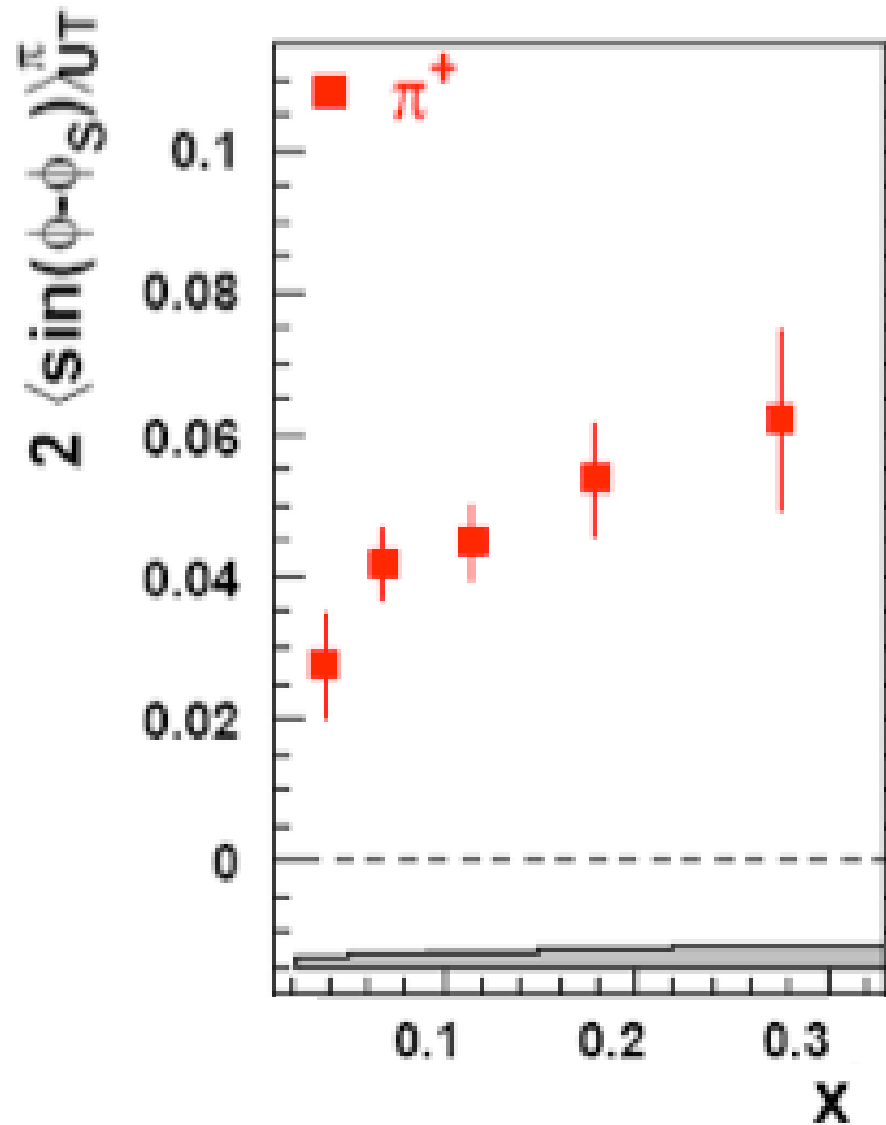
- Brief introduction
 - SSA theory
 - Sign change
- Sivers function and current predictions
 - Global fitting from SIDIS
 - Predictions for DY
 - Some other channels
- Consequence of DY measurements

Experimental data on single transverse spin asymmetry

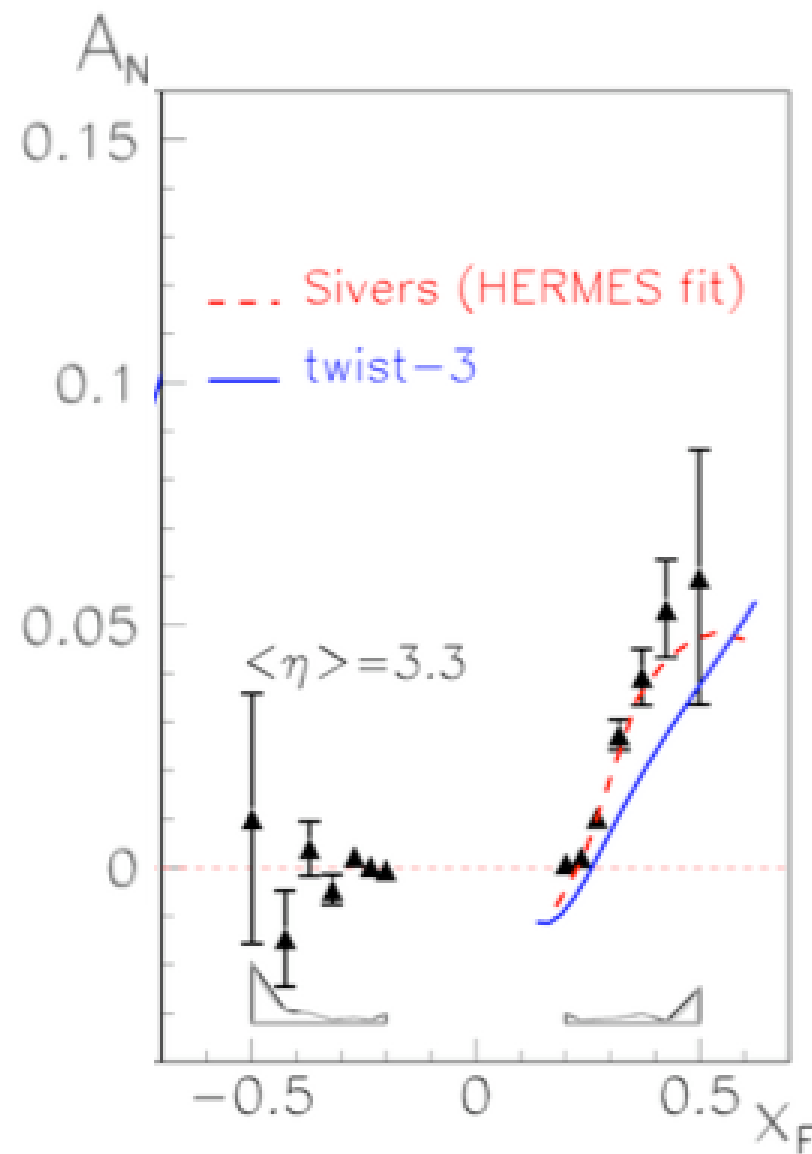
- Single transverse spin asymmetries (SSAs) have been observed in various experiments at different CM energies

- PHENIX, BRAHMS, COMPASS, JLAB, too

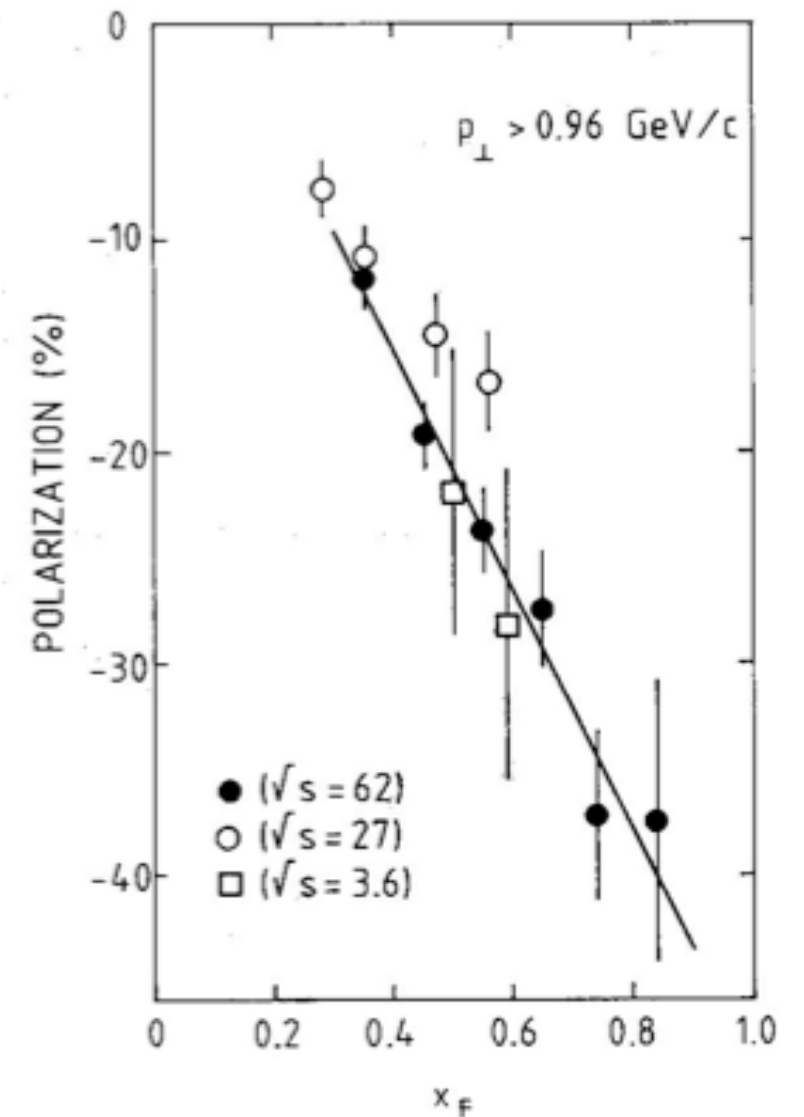
$$A_N \equiv \frac{\Delta\sigma(\ell, \vec{s})}{\sigma(\ell)} = \frac{\sigma(\ell, \vec{s}) - \sigma(\ell, -\vec{s})}{\sigma(\ell, \vec{s}) + \sigma(\ell, -\vec{s})}$$



$\ell + p^\uparrow \rightarrow \ell' + \pi + X$
HERMES



$p + p^\uparrow \rightarrow \pi + X$
STAR



$p + p \rightarrow \Lambda^\uparrow + X$
E704

SSA vanishes at leading twist in collinear factorization

Kane, Pumplin, Repko, 1978

- At leading twist formalism: partons are collinear

$$\sigma(s_T) \sim \left| \begin{array}{c} \text{Diagram (a)} \\ \text{Diagram (b)} \\ \vdots \end{array} \right|^2 \rightarrow \Delta\sigma(s_T) \sim \text{Re}[(a)] \cdot \text{Im}[(b)]$$

The diagram shows two Feynman diagrams, (a) and (b), representing parton interactions. Diagram (a) shows a parton with momentum p and spin s_p interacting with a parton with momentum k . Diagram (b) shows a similar interaction but with a different internal structure. The diagrams are summed and squared to give the cross-section $\sigma(s_T)$. The difference in cross-sections $\Delta\sigma(s_T)$ is proportional to the real part of (a) times the imaginary part of (b).

- generate phase from loop diagrams, proportional to α_s
- helicity is conserved for massless partons, helicity-flip is proportional to current quark mass m_q

Therefore we have

$$A_N \sim \alpha_s \frac{m_q}{\sqrt{s}} \rightarrow 0$$

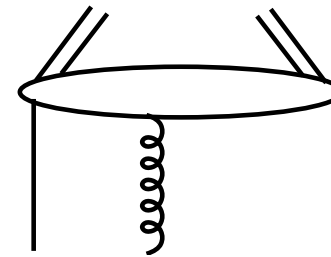
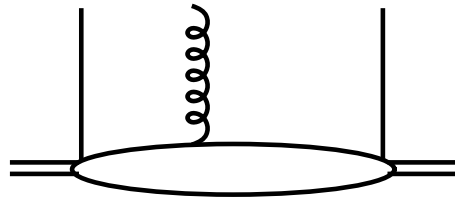
- $A_N \neq 0$: result of parton's transverse motion or correlations!

Two mechanisms to generate SSA in QCD

Collinear twist-3 factorization approach

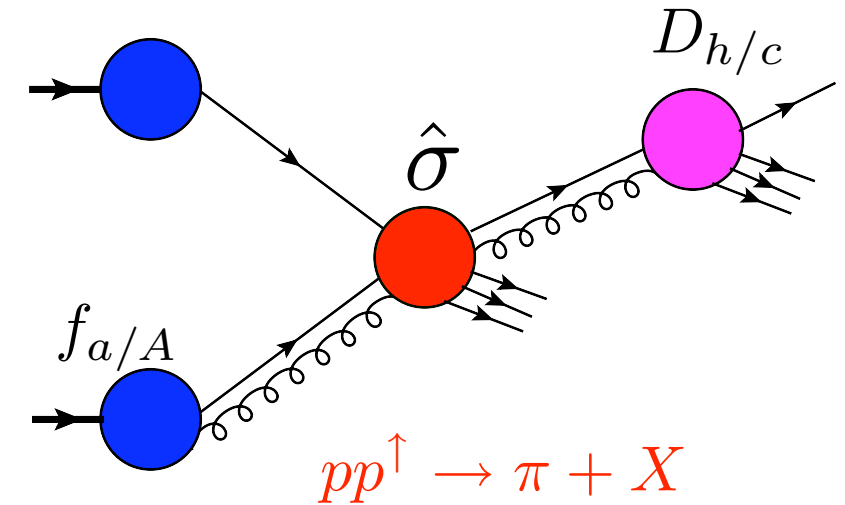
$$\sigma(p_h, s_\perp) \propto f_{a/A}^{s_\perp}(x) \otimes D_{h/c}(z) \otimes \hat{\sigma}_{\text{parton}}$$

- Twist-3 three-parton correlation functions (PDFs)
- Twist-3 three-parton fragmentation functions



Efremov-Teryaev 82, 84, Qiu-Sterman 91, 98, ...

Koike, 02, Kang, Yuan, Zhou 2010



TMD approach: Transverse Momentum Dependent distributions probe the parton's intrinsic transverse momentum

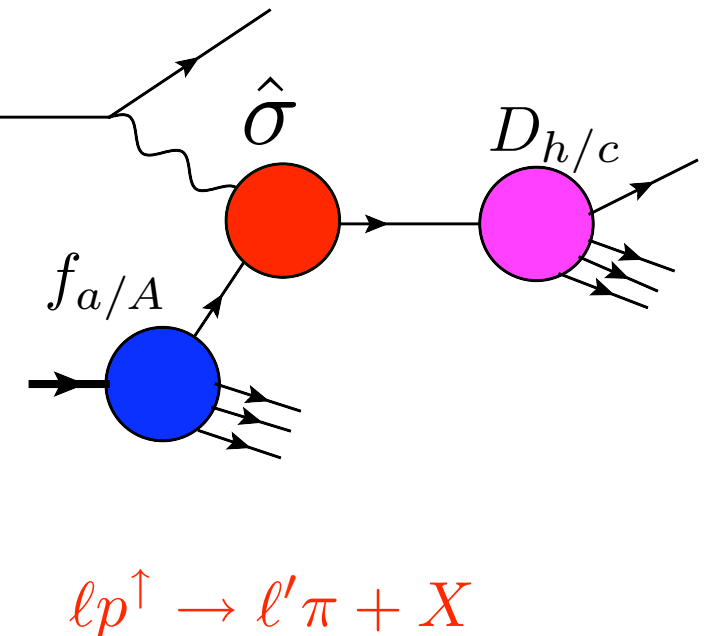
$$\sigma(p_h, s_\perp) \propto f_{a/A}(x, k_\perp) \otimes D_{h/c}(z, p_\perp) \otimes \hat{\sigma}_{\text{parton}}$$

- Sivers function: in Parton Distribution Function (PDF)

Sivers 90

- Collins function: in Fragmentation Function (FF)

Collins 93



Relation between twist-3 and TMD approaches

- They apply in different kinematic domain:

- TMD approach: need TMD factorization, applies for the process with two observed momentum scales: DY at small Q_T

$$Q_1 \gg Q_2 \left\{ \begin{array}{l} Q_1 \text{ necessary for pQCD factorization to have a chance} \\ Q_2 \text{ sensitive to parton's transverse momentum} \end{array} \right.$$

- Collinear factorization approach: more relevant for single scale hard process: inclusive pion production at pp collision

- They generate same results in the overlap region when they both apply:

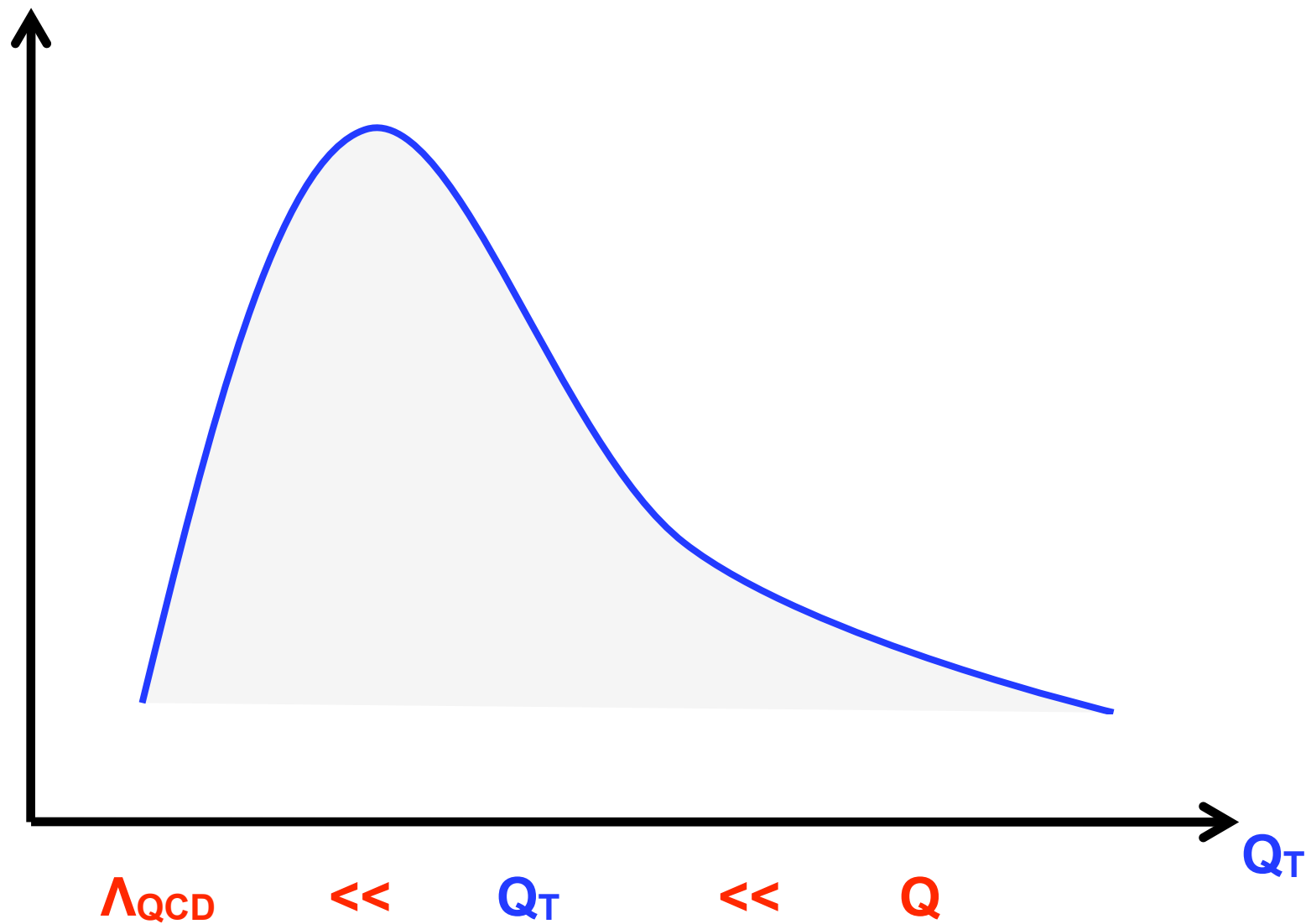
- Twist-3 three-parton correlation in distribution \longleftrightarrow Siverson function

Ji, Qiu, Vogelsang, Yuan, 2006, ...

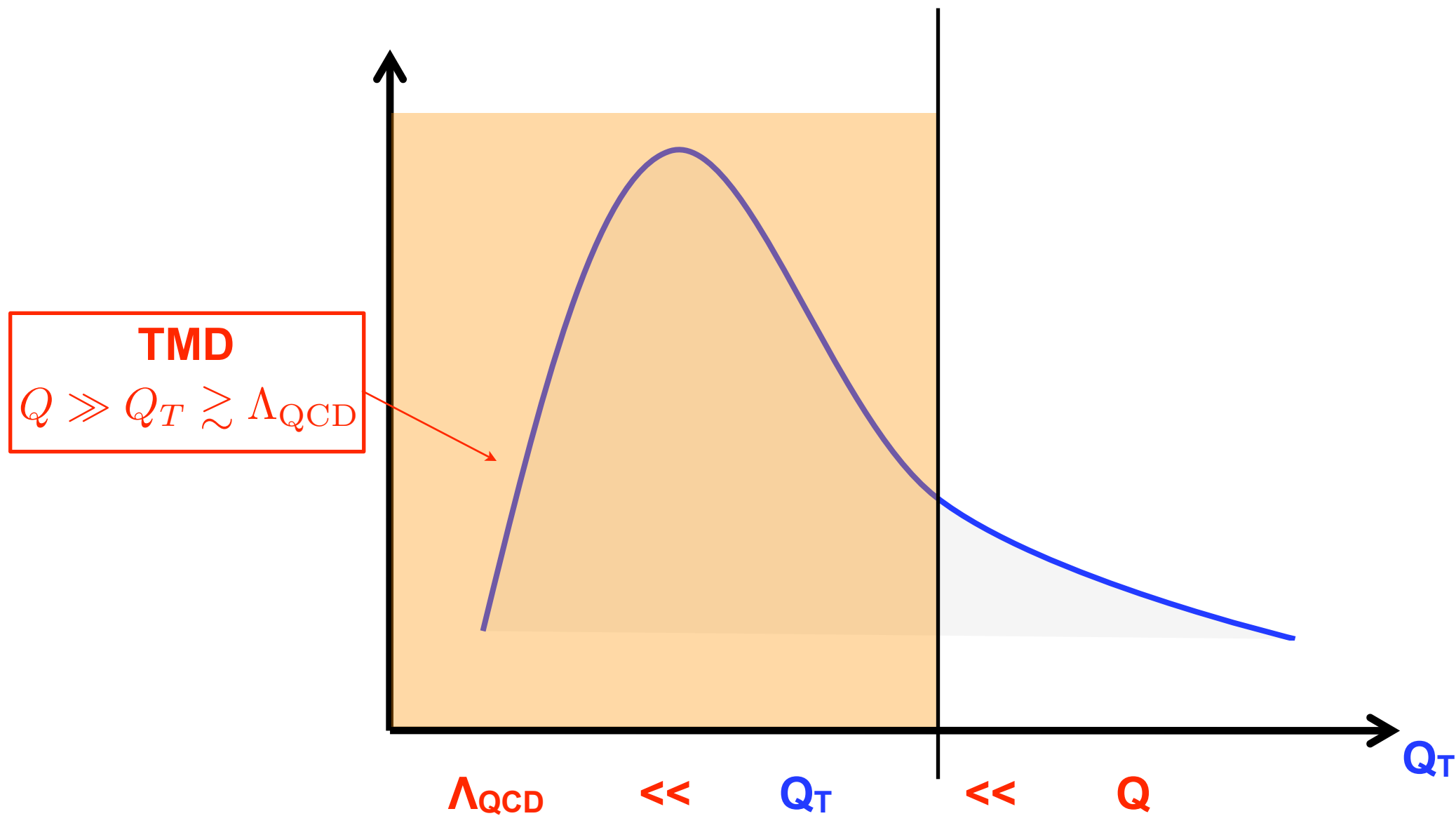
- Twist-3 three-parton correlation in fragmentation \longleftrightarrow Collins function

Zhou, Yuan, 2009, Kang, Yuan, Zhou, 2010

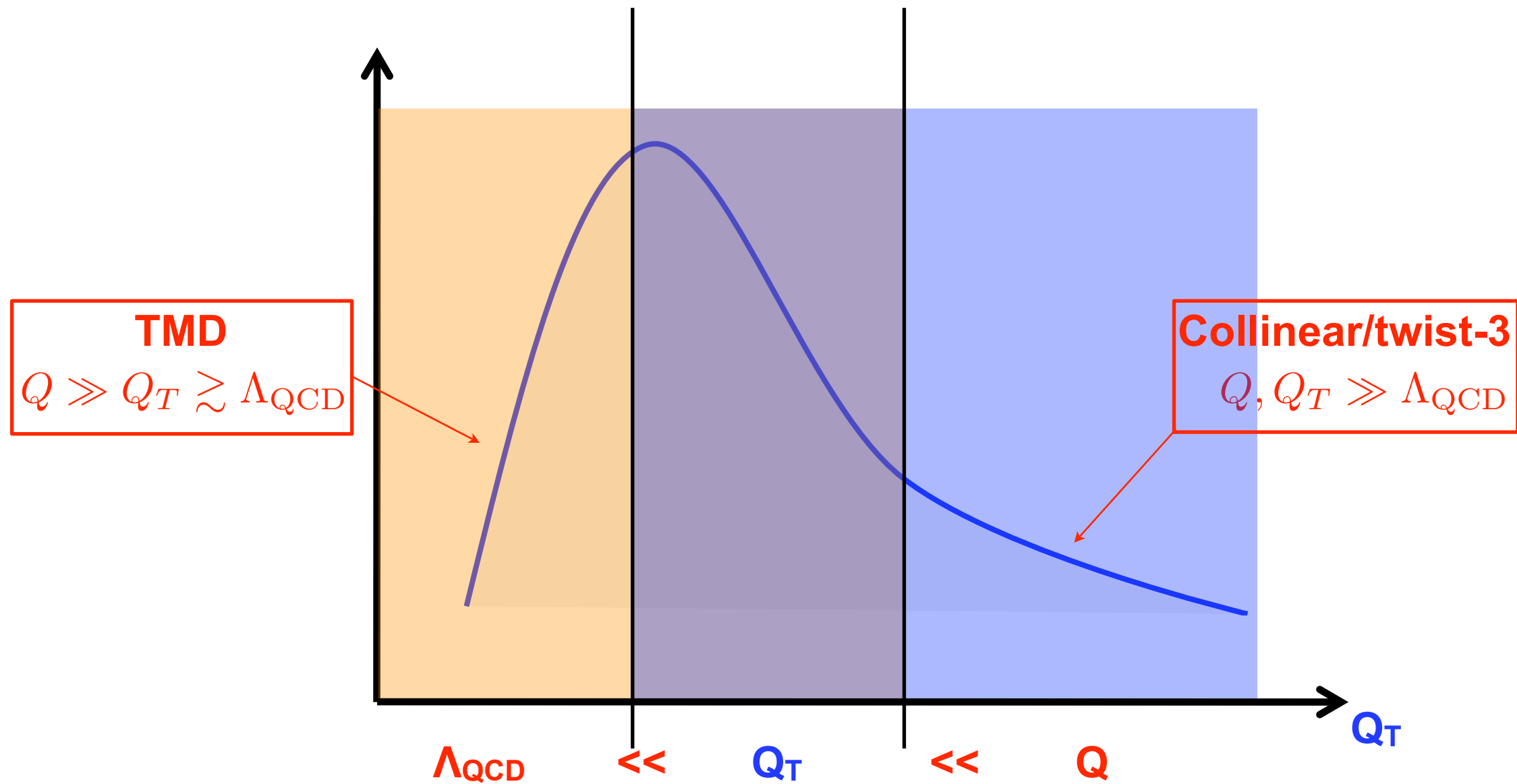
A unified picture for Drell-Yan (leading Q_T/Q)



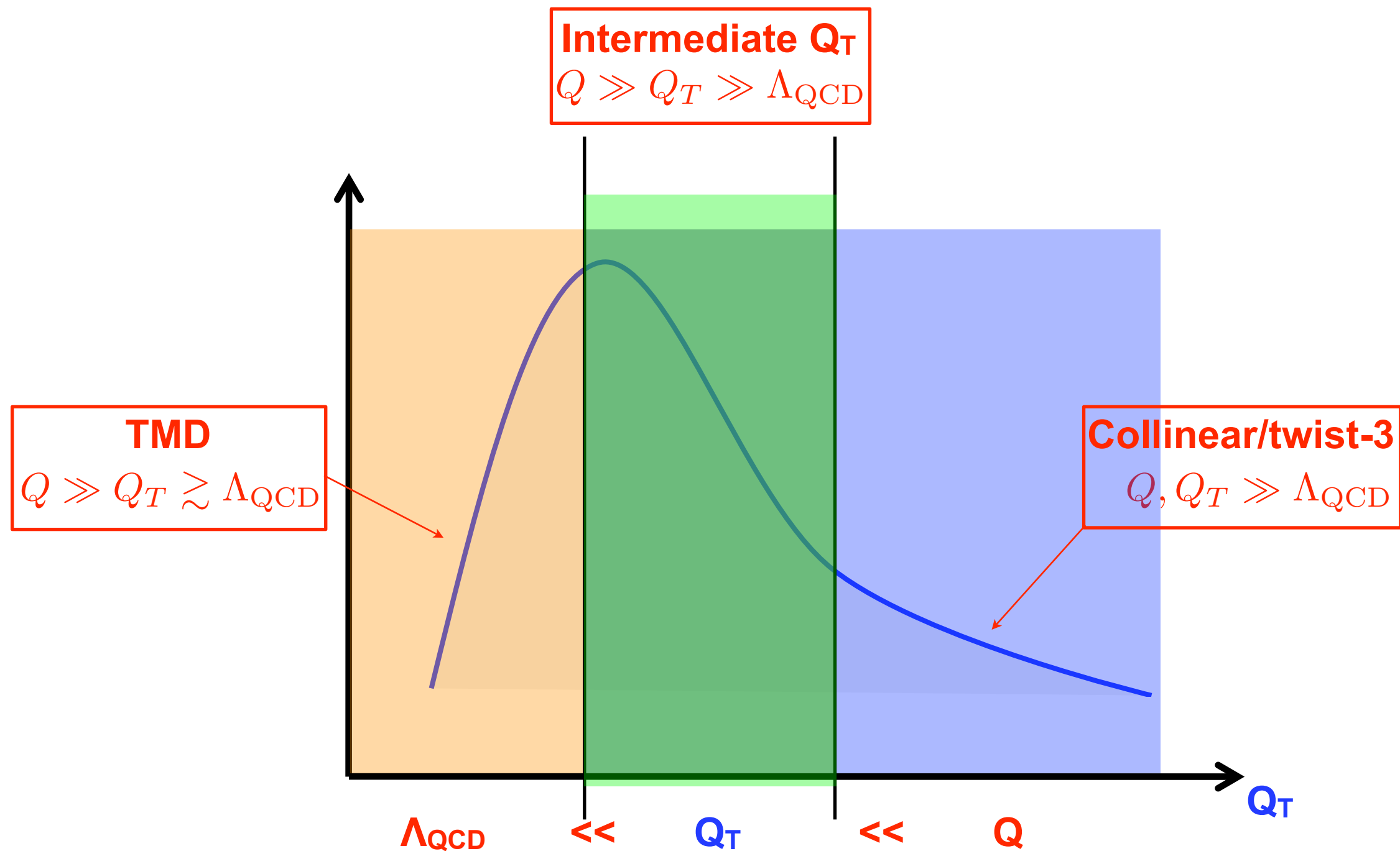
A unified picture for Drell-Yan (leading Q_T/Q)



A unified picture for Drell-Yan (leading Q_T/Q)



A unified picture for Drell-Yan (leading Q_T/Q)





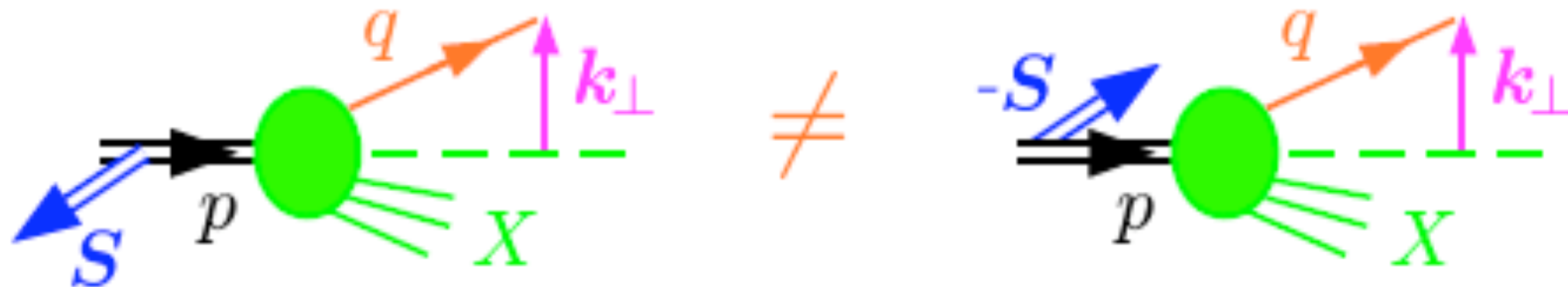
Major difference in these two approaches

- Collinear factorization approach:
 - All the twist-3 correlation functions (both in distribution and fragmentation side) are universal
 - Any process-dependent part is in the hard-part, which is calculable
- However, the TMD function in TMD approach MIGHT not be universal
 - Sivers function is NOT universal
Collins 02, Boer, Mulders, Pijlman, 03, Collins, Metz, 04, Kang, Qiu, 09, ...
 - Collins function is universal
Metz 02, Collins, Metz, 04, Yuan, 08, Gamberg, Mukerjee, Mulders, 08, Meissner, Metz, 08, Zhou, Yuan, 09, Boer, Kang, Vogelsang, Yuan, 10...

TMD approach: Sivers function

- An asymmetric parton distribution in a polarized hadron (k_t correlated with the spin of the hadron)

$$f_{q/h^\uparrow}(x, \mathbf{k}_\perp, \vec{S}) \equiv \underbrace{f_{q/h}(x, k_\perp)}_{\text{Spin-independent}} + \underbrace{\frac{1}{2} \Delta^N f_{q/h^\uparrow}(x, k_\perp)}_{\text{Spin-dependent}} \vec{S} \cdot \hat{p} \times \hat{\mathbf{k}}_\perp$$



From experiments to theory: QCD kt-factorization

- One measures cross sections in the experiment, and then use theory to connect to the relevant distributions (hadron structure)

- SIDIS

$$\sigma \sim \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$
$$\approx \text{[diagram 4]} \otimes \text{[diagram 5]}$$

PDFs with SIDIS gauge link

The diagrams illustrate the perturbative expansion of the SIDIS cross section. The first row shows the cross section σ as a sum of diagrams: a leading-order diagram with a quark line and a gluon exchange, followed by higher-order diagrams with additional gluon exchanges. The second row shows the factorization of the cross section into a hard scattering kernel (represented by a diagram with a cross in a circle) and a parton distribution function (represented by a diagram with a cross in a circle). The text "PDFs with SIDIS gauge link" indicates that the PDFs are defined with a specific gauge link configuration.

- DY

$$\sigma \sim \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$
$$\approx \text{[diagram 4]} \otimes \text{[diagram 5]}$$

PDFs with DY gauge link

The diagrams illustrate the perturbative expansion of the DY cross section. The first row shows the cross section σ as a sum of diagrams: a leading-order diagram with a quark line and a gluon exchange, followed by higher-order diagrams with additional gluon exchanges. The second row shows the factorization of the cross section into a hard scattering kernel (represented by a diagram with a cross in a circle) and a parton distribution function (represented by a diagram with a cross in a circle). The text "PDFs with DY gauge link" indicates that the PDFs are defined with a specific gauge link configuration.

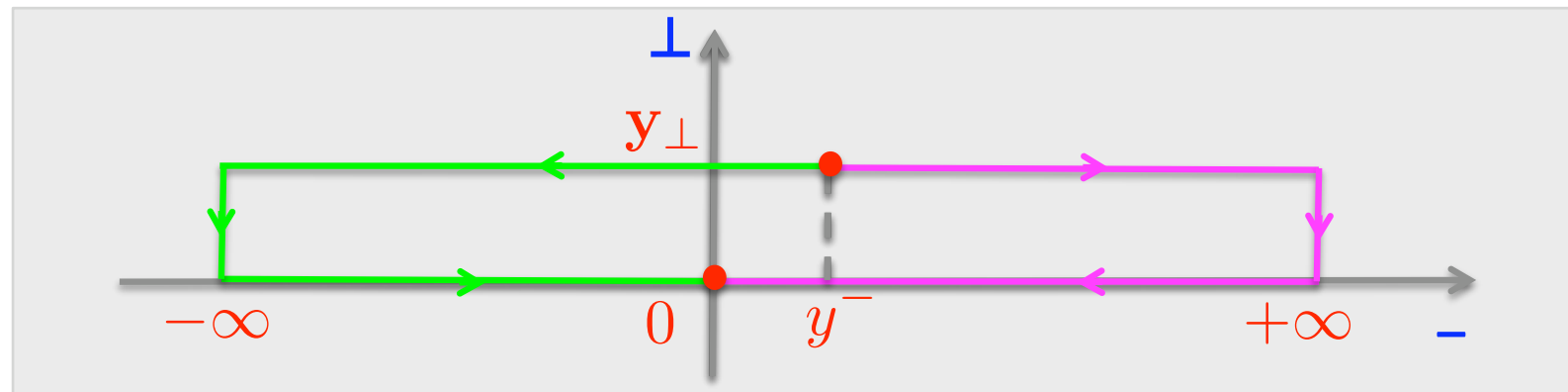
Non-universality of the Siverson function

- Different gauge link for gauge-invariant TMD distribution in SIDIS and DY

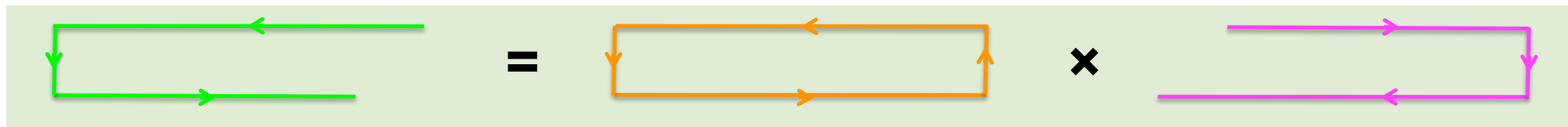
$$f_{q/h^\uparrow}(x, \mathbf{k}_\perp, \vec{S}) = \int \frac{dy^- d^2 y_\perp}{(2\pi)^3} e^{ixp^+ y^- - i \mathbf{k}_\perp \cdot \mathbf{y}_\perp} \langle p, \vec{S} | \bar{\psi}(0^-, \mathbf{0}_\perp) \boxed{\text{Gauge link}} \frac{\gamma^+}{2} \psi(y^-, \mathbf{y}_\perp) | p, \vec{S} \rangle$$

- **SIDIS:** $\Phi_n^\dagger(\{+\infty, 0\}, \mathbf{0}_\perp) \Phi_{n_\perp}^\dagger(+\infty, \{\mathbf{y}_\perp, \mathbf{0}_\perp\}) \Phi_n(\{+\infty, y^-\}, \mathbf{y}_\perp)$

- **DY:** $\Phi_n^\dagger(\{-\infty, 0\}, \mathbf{0}_\perp) \Phi_{n_\perp}^\dagger(-\infty, \{\mathbf{y}_\perp, \mathbf{0}_\perp\}) \Phi_n(\{-\infty, y^-\}, \mathbf{y}_\perp)$



Wilson Loop $\sim \exp \left[-ig \int_\Sigma d\sigma^{\mu\nu} F_{\mu\nu} \right]$ Area is NOT zero



- For a fixed spin state:

$$f_{q/h^\uparrow}^{\text{SIDIS}}(x, \mathbf{k}_\perp, \vec{S}) \neq f_{q/h^\uparrow}^{\text{DY}}(x, \mathbf{k}_\perp, \vec{S})$$

Time-reversal modified universality of the Sivers function

- Relation between Sivers functions in SIDIS and DY

- From P and T invariance:

$$f_{q/h\uparrow}^{\text{SIDIS}}(x, \mathbf{k}_{\perp}, \vec{S}) = f_{q/h\uparrow}^{\text{DY}}(x, \mathbf{k}_{\perp}, -\vec{S})$$

- **Spin-averaged parton distribution function is universal**

$$f_{q/h}(x, k_{\perp}) = \frac{1}{2} \left[f_{q/h\uparrow}(x, \mathbf{k}_{\perp}, \vec{S}) + f_{q/h\uparrow}(x, \mathbf{k}_{\perp}, -\vec{S}) \right]$$

- From the definition of Sivers function:

$$\Delta^N f_{q/h\uparrow}(x, k_{\perp}) \vec{S} \cdot \hat{p} \times \hat{\mathbf{k}}_{\perp} = f_{q/h\uparrow}(x, \mathbf{k}_{\perp}, \vec{S}) - f_{q/h\uparrow}(x, \mathbf{k}_{\perp}, -\vec{S})$$

- One can derive:

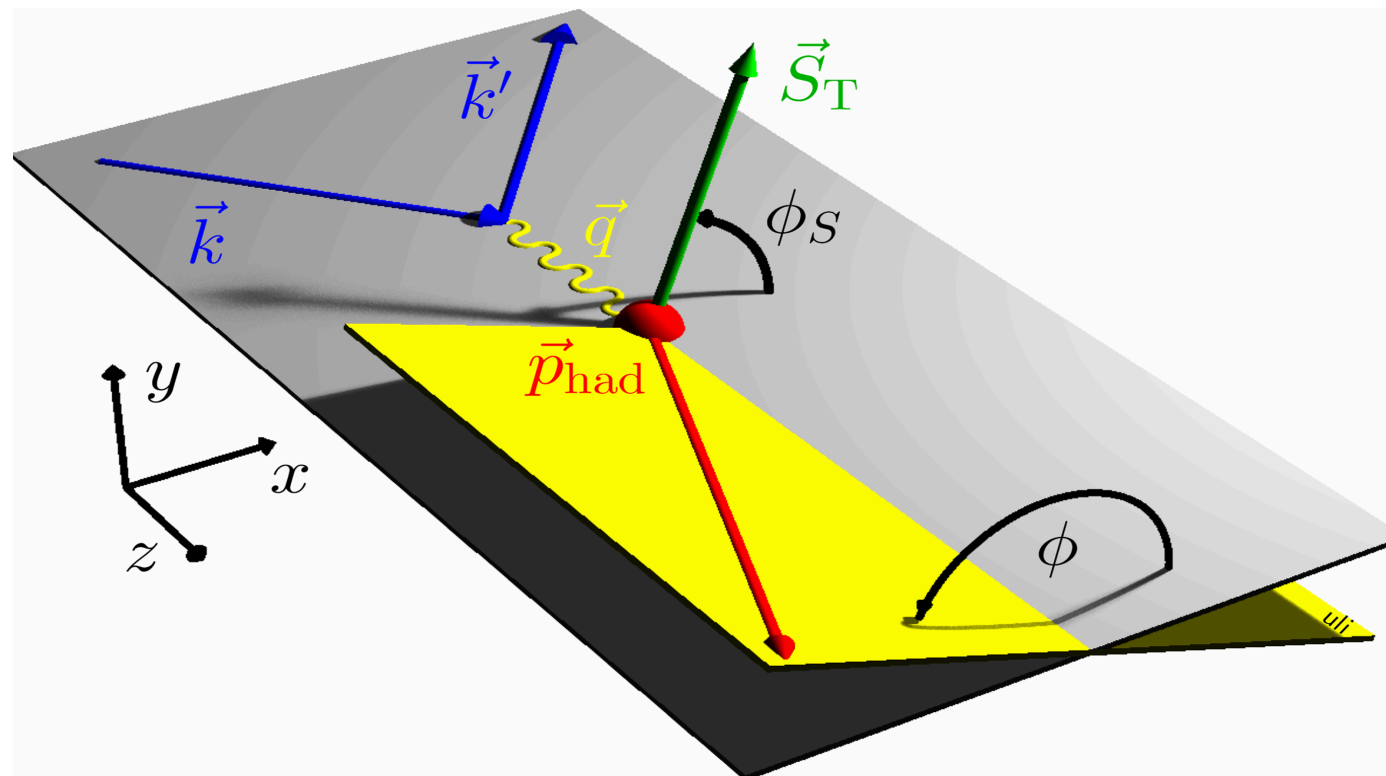
$$\Delta^N f_{q/h\uparrow}^{\text{SIDIS}}(x, k_{\perp}) = -\Delta^N f_{q/h\uparrow}^{\text{DY}}(x, k_{\perp})$$

Most critical test for TMD approach to SSA

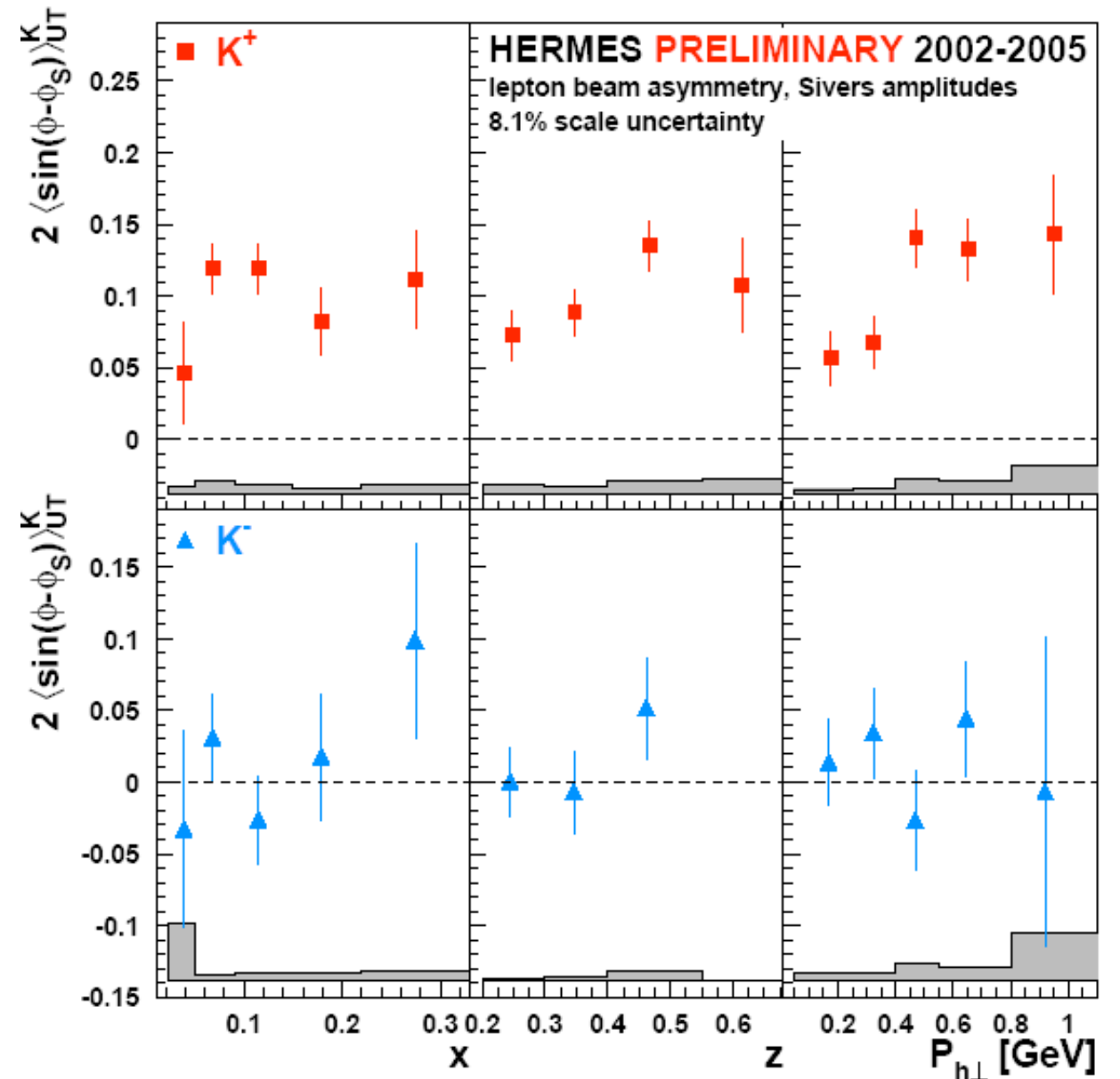
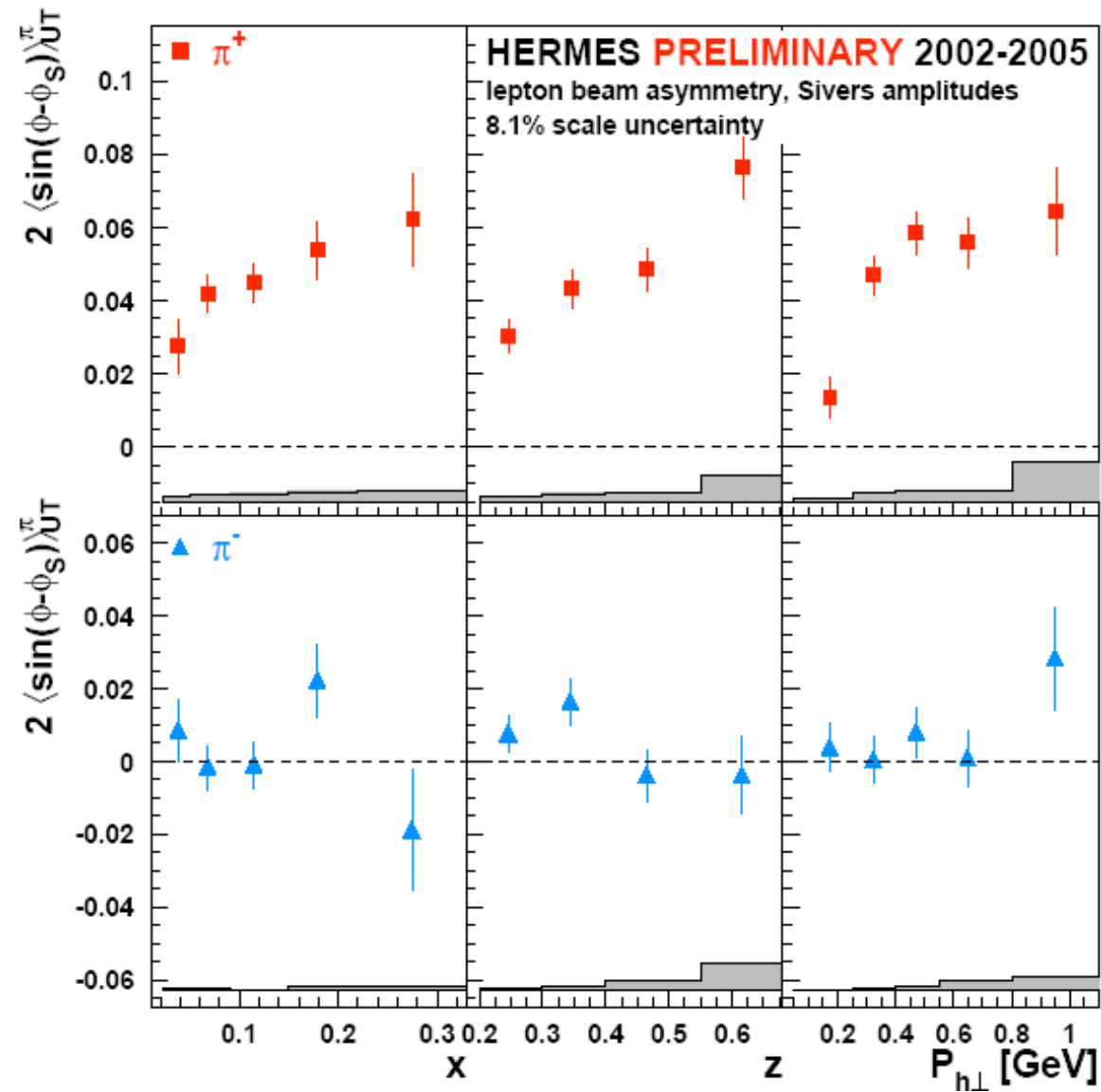
Current Sivers function from SIDIS

- Sivers and Collins can be separately extracted from SIDIS

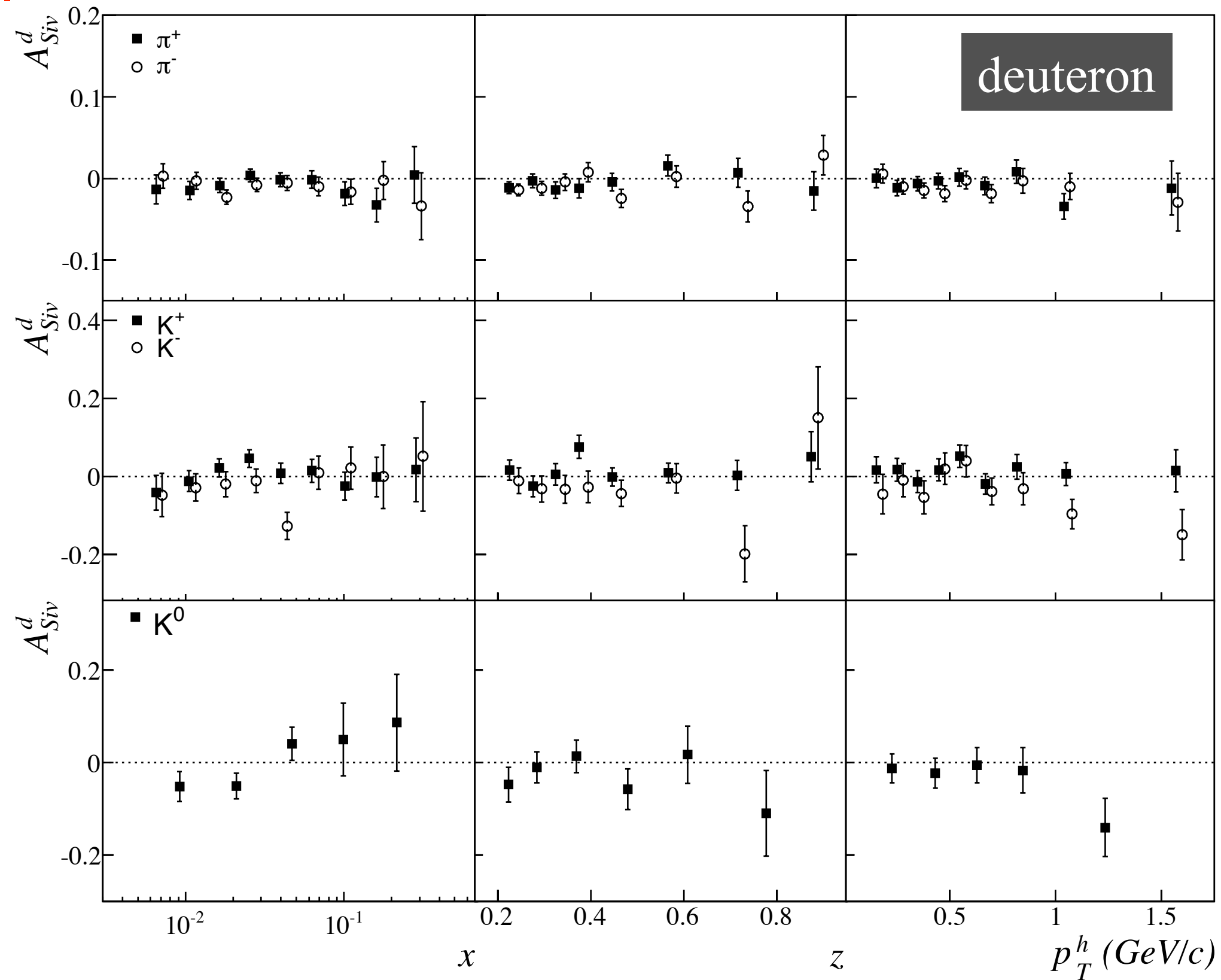
$$\Delta\sigma \propto A_{\text{UT}}^{\text{Collins}} \sin(\phi + \phi_S) + A_{\text{UT}}^{\text{Sivers}} \sin(\phi - \phi_S)$$



HERMES: Preliminary results on Proton (NOT zero)



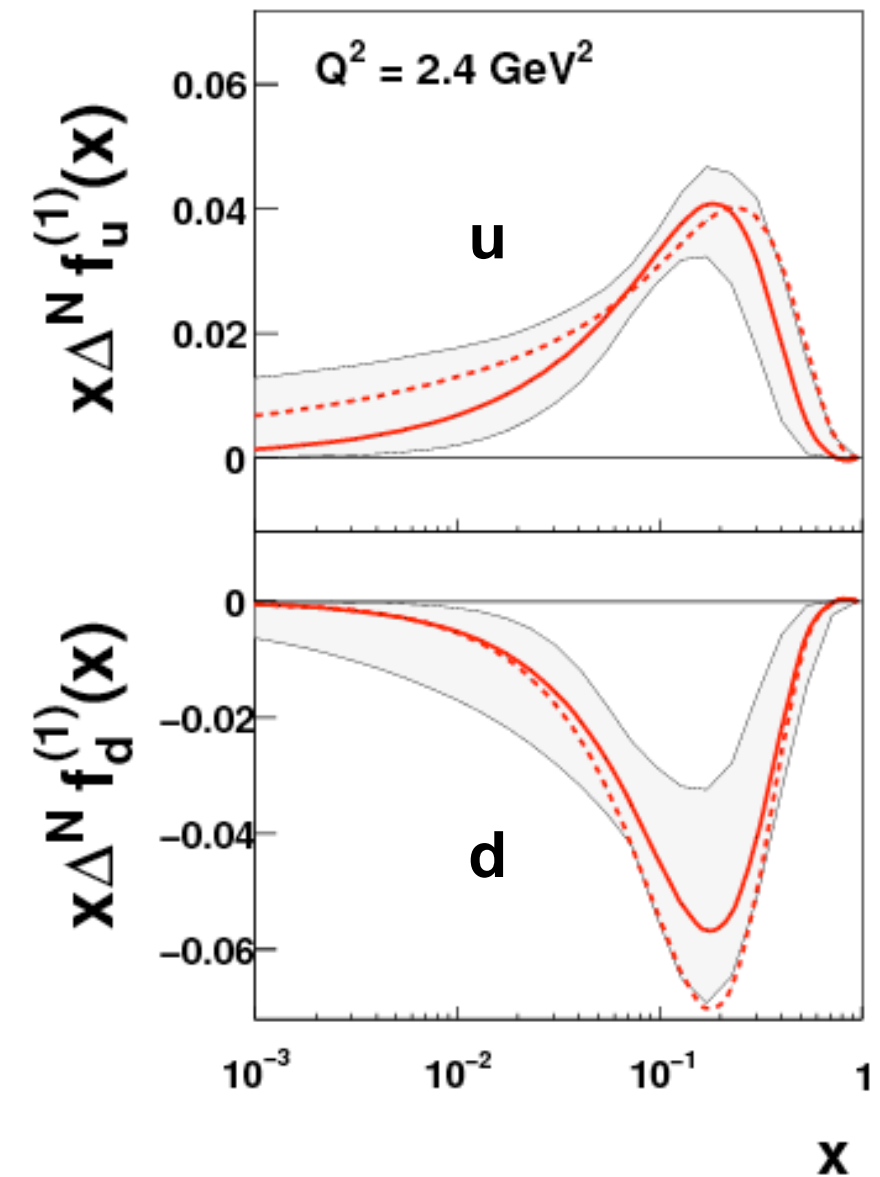
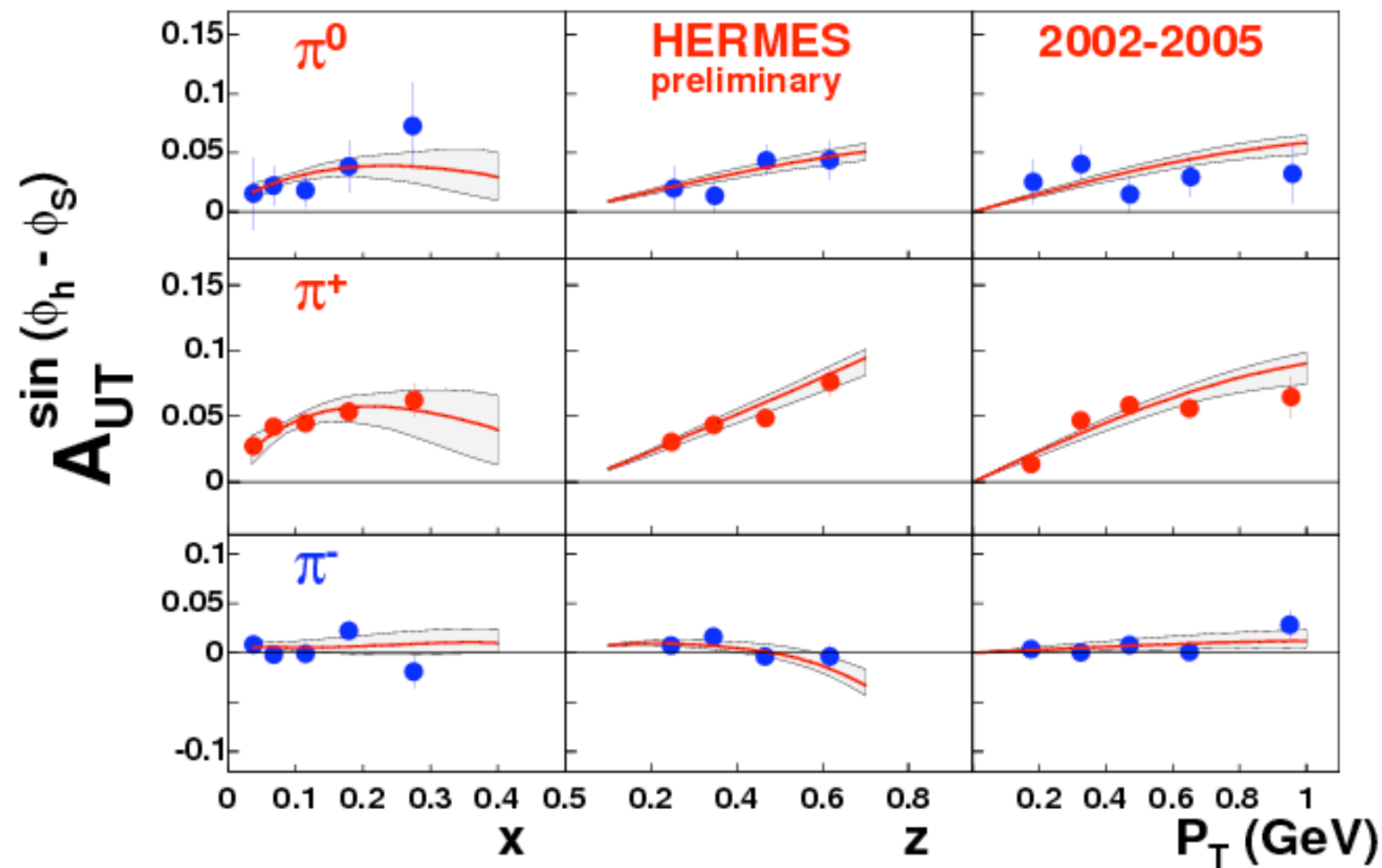
COMPASS: Deuteron target (small or zero)



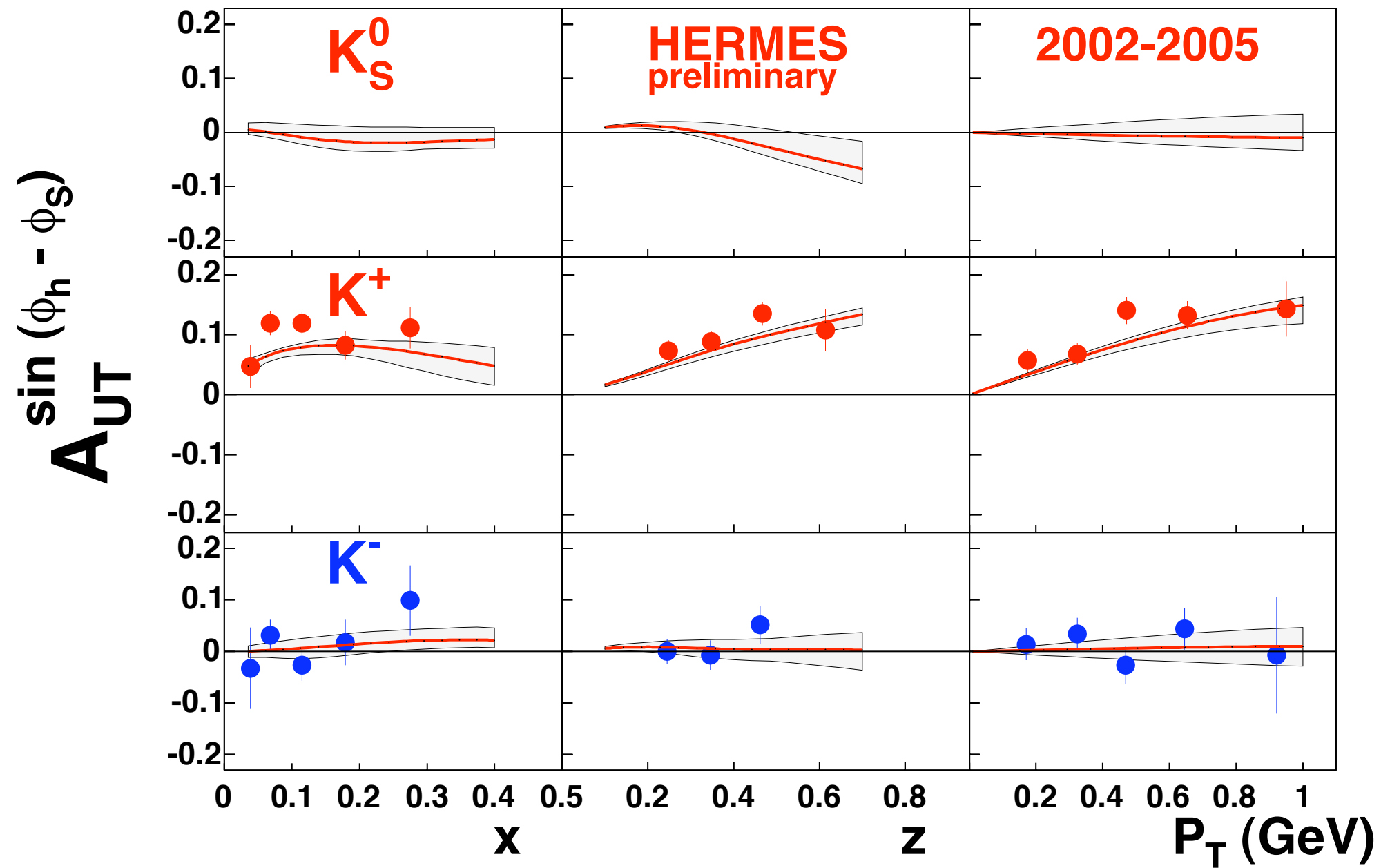
Sivers function from SIDIS: Current Global Analysis

- Includes HERMES Proton data and COMPASS Deuteron data

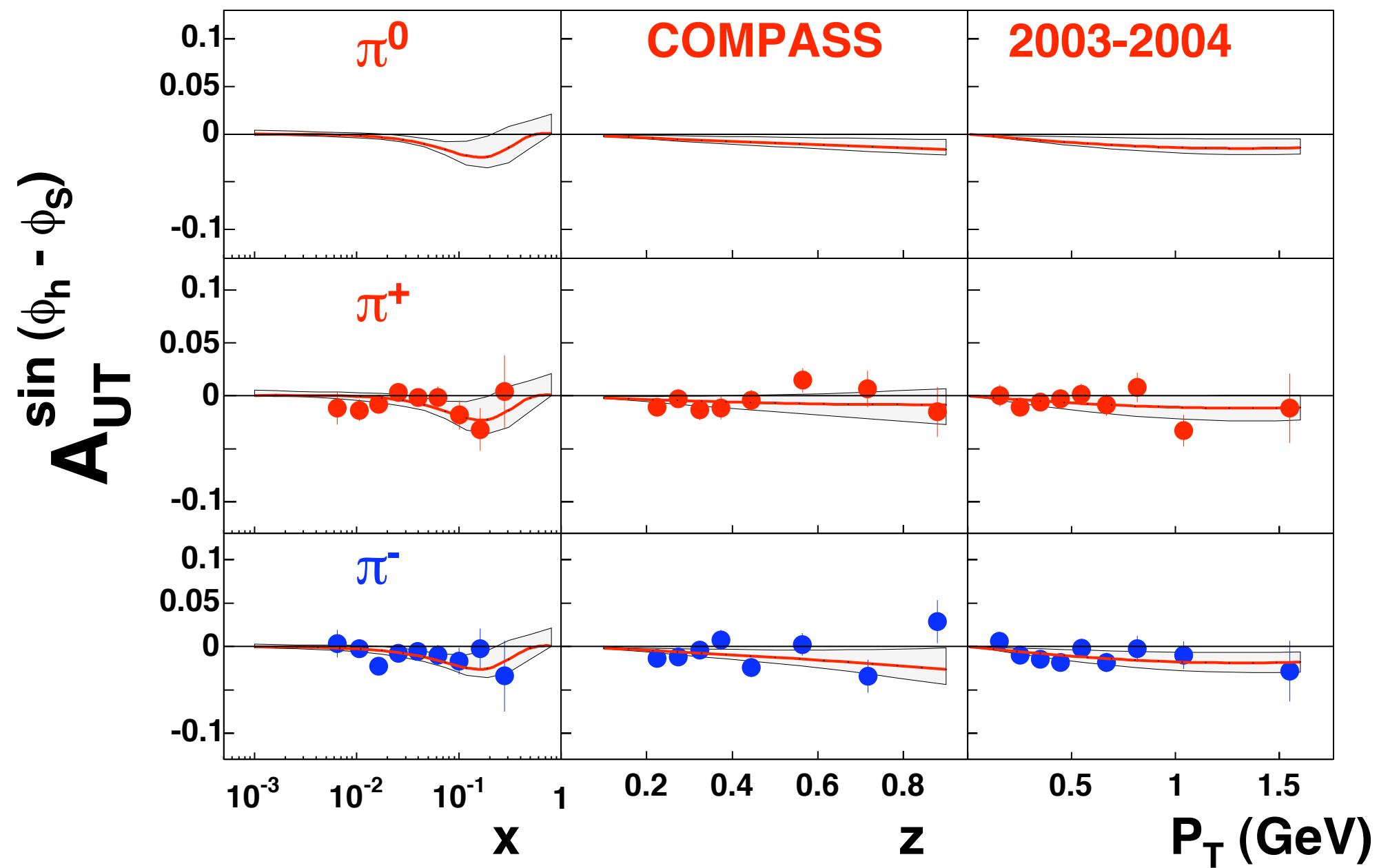
Anselmino, et.al., 2009



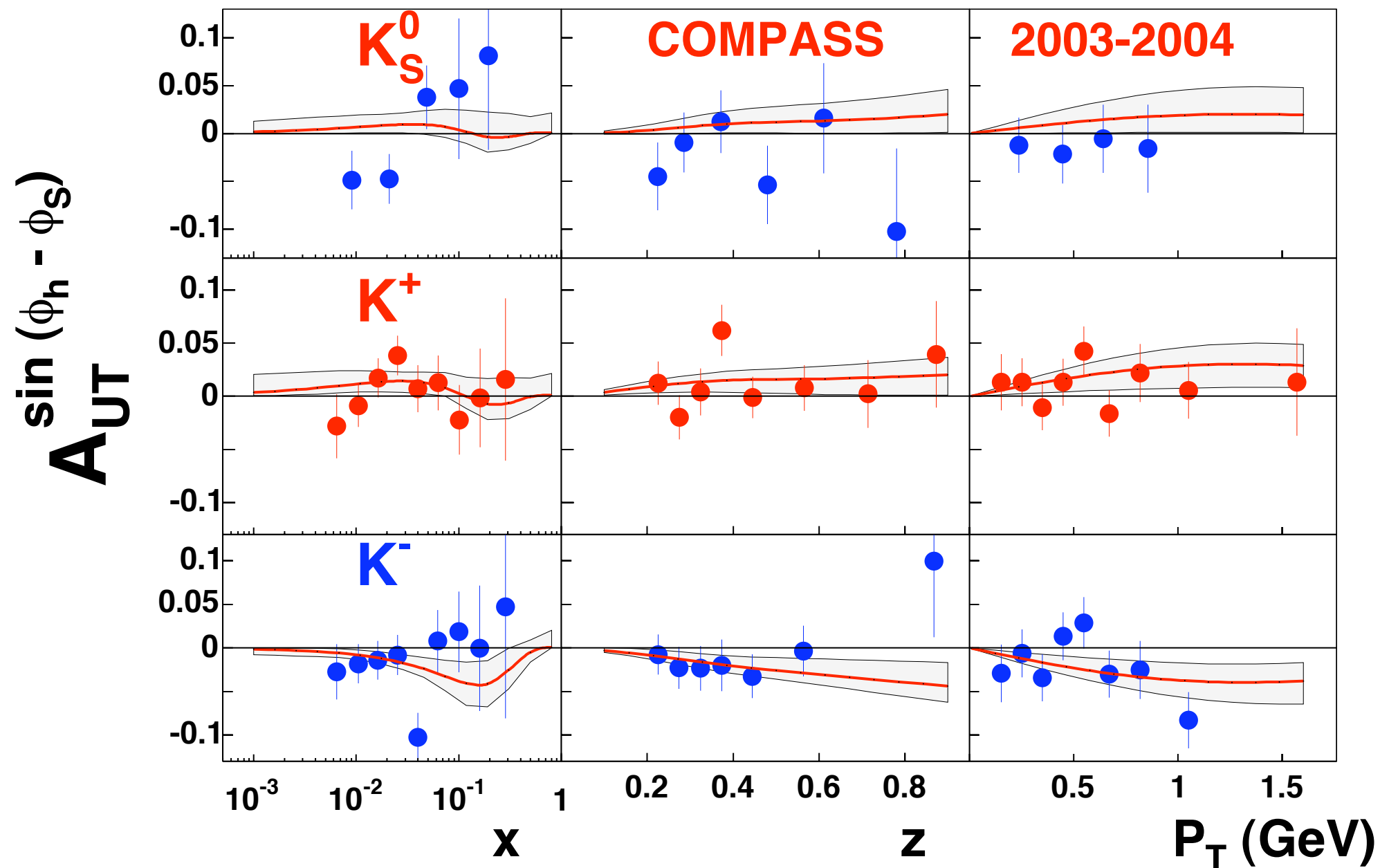
Comparison with HERMES Proton: Kaons



Comparison with COMPASS Deuteron: Pions

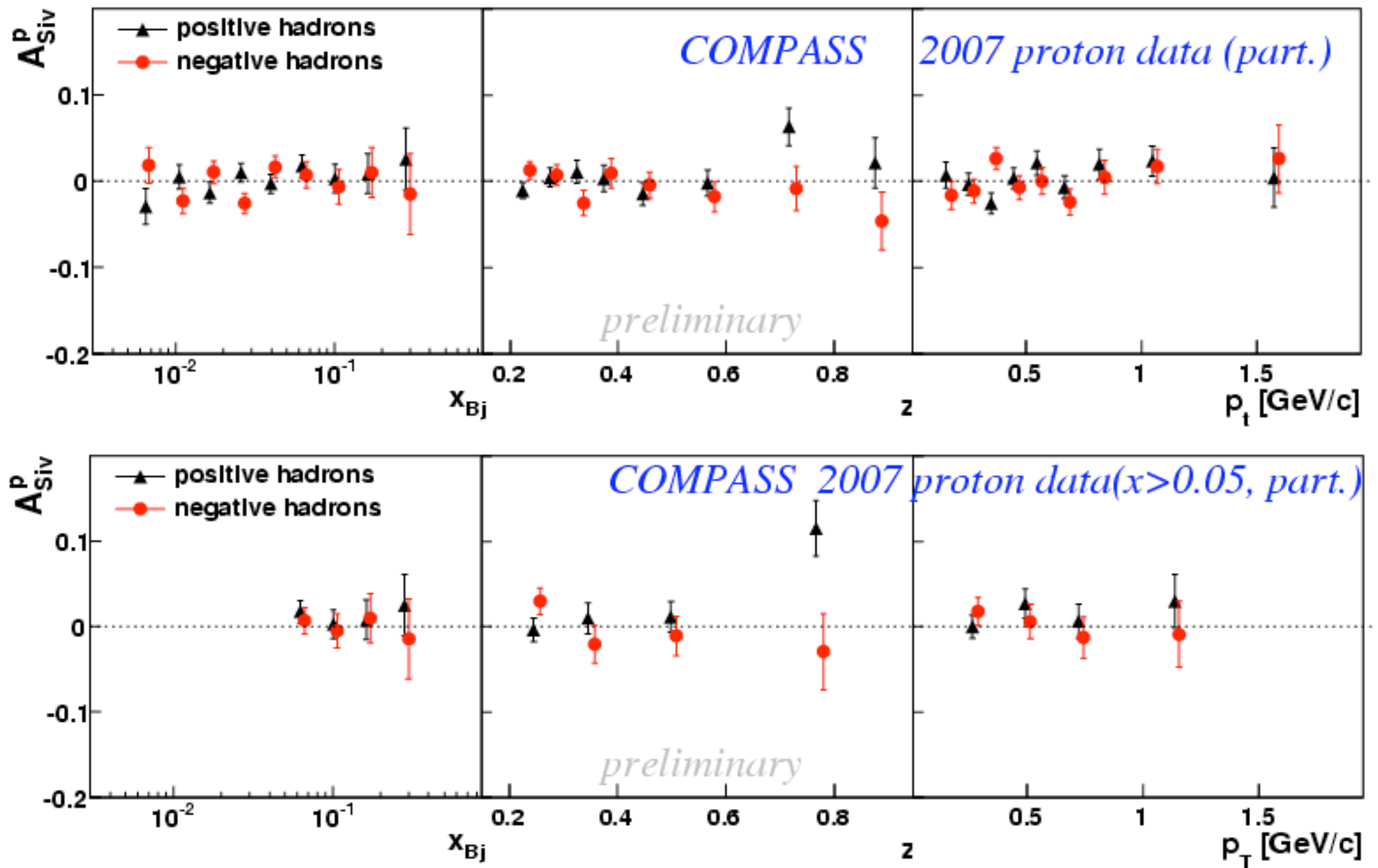


Comparison with COMPASS Deuteron: Kaons



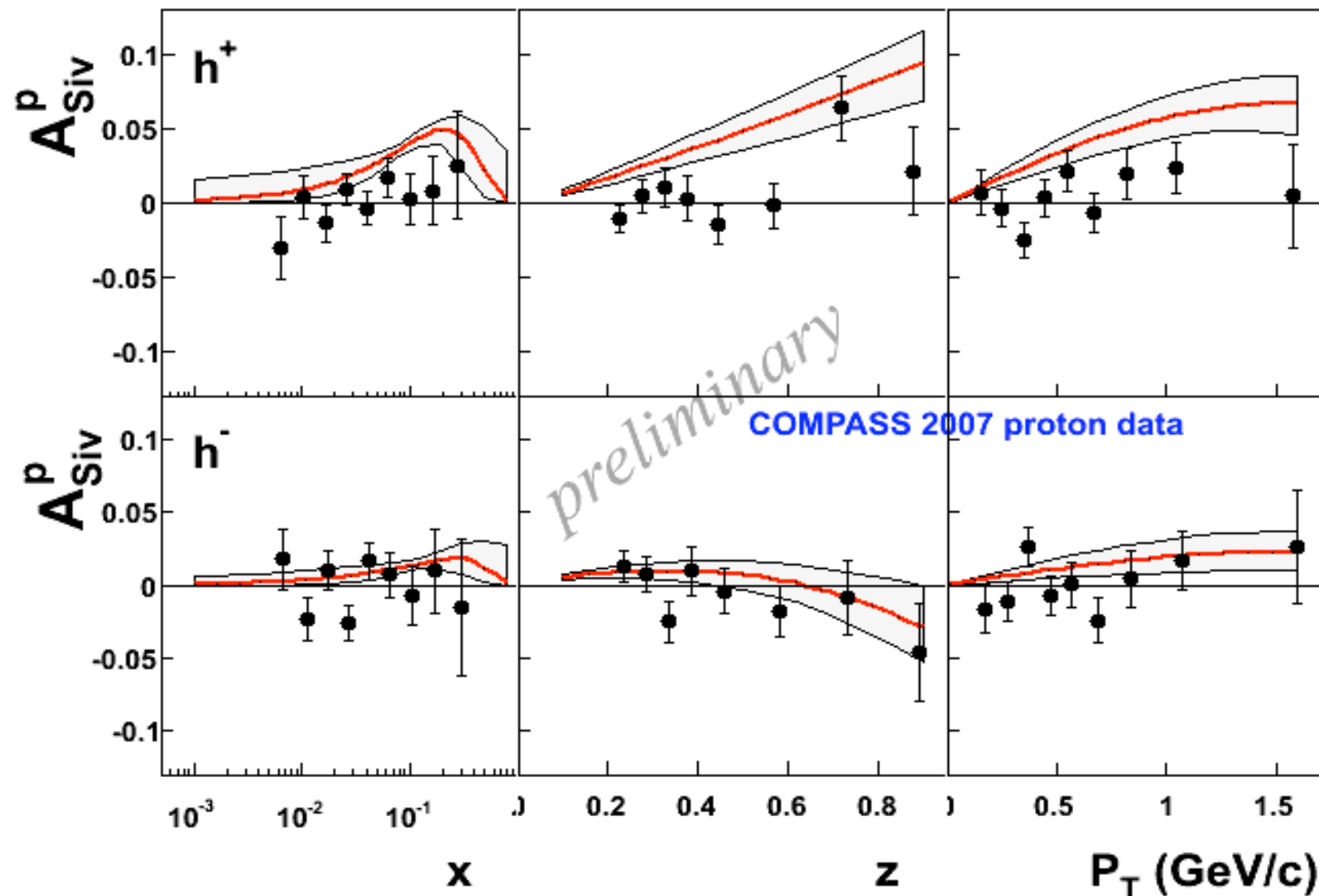
What about COMPASS proton?

COMPASS: proton (small or zero)



COMPASS Proton compare with theory

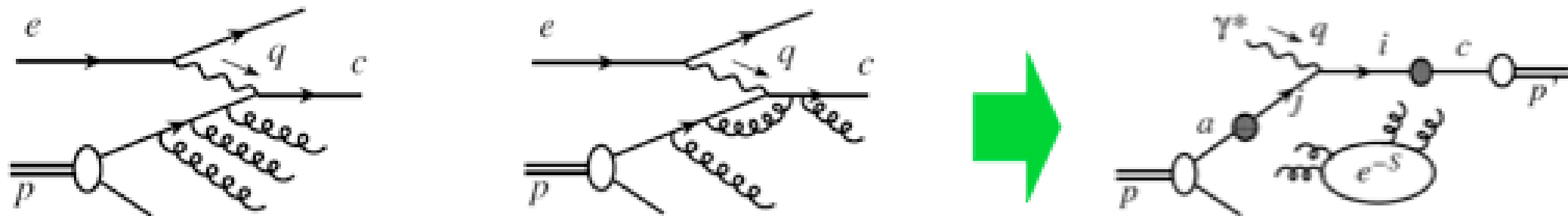
- The predictions do not seem to be consistent with COMPASS proton



- Incompatible between HERMES and COMPASS data on Sivers effect
- Q^2 makes a difference: $\langle Q^2 \rangle \sim 10$ (COMPASS) and 2.4 (HERMES) GeV^2
- QCD resummation?

QCD resummation: Sudakov suppression at high Q^2

- QCD resummation

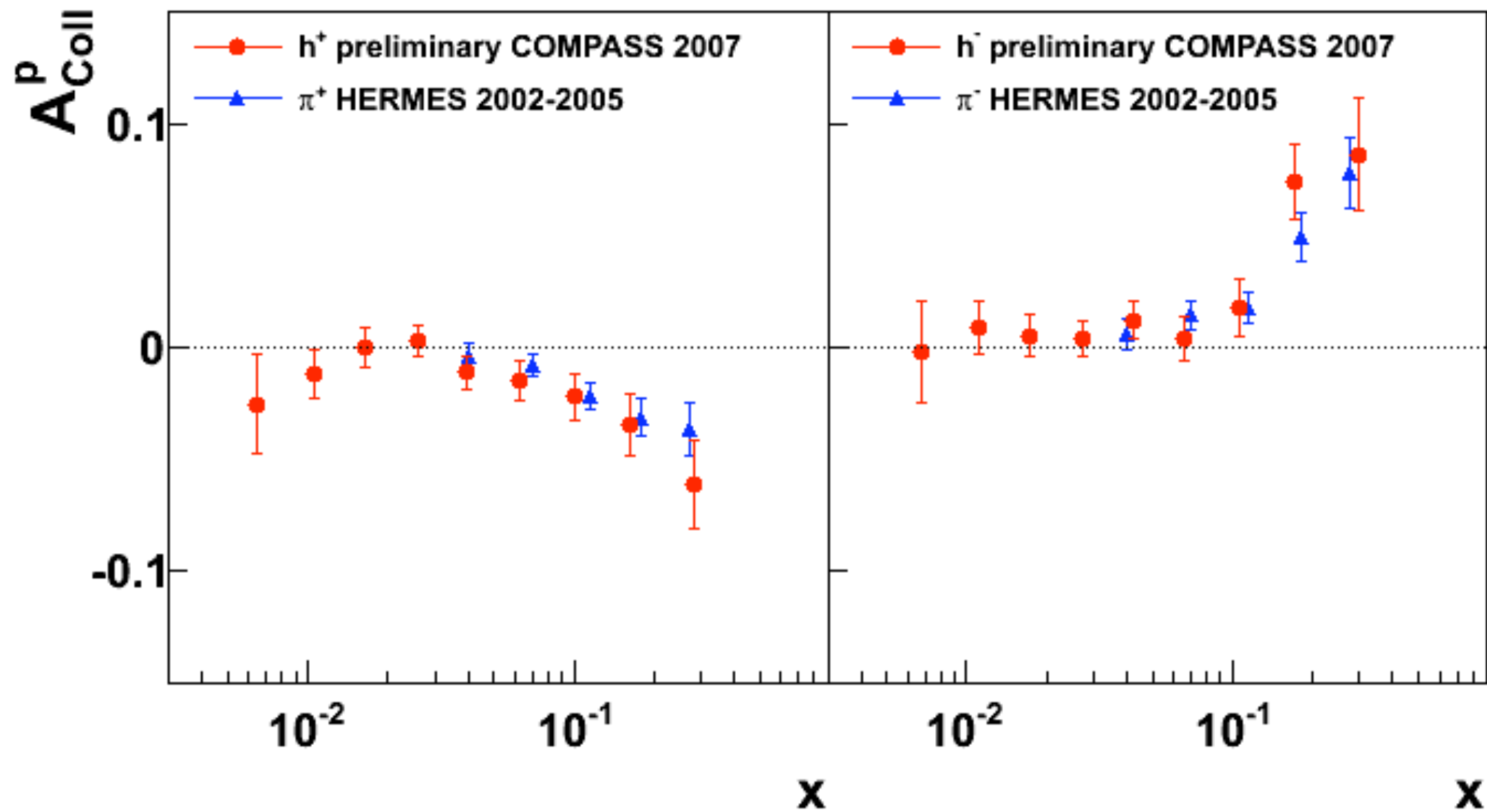


initial-state and final-state soft gluon radiations generate

large logarithms: $\frac{1}{q_T^2} \alpha_s^n \log^{2n-1}(Q^2/q_T^2)$

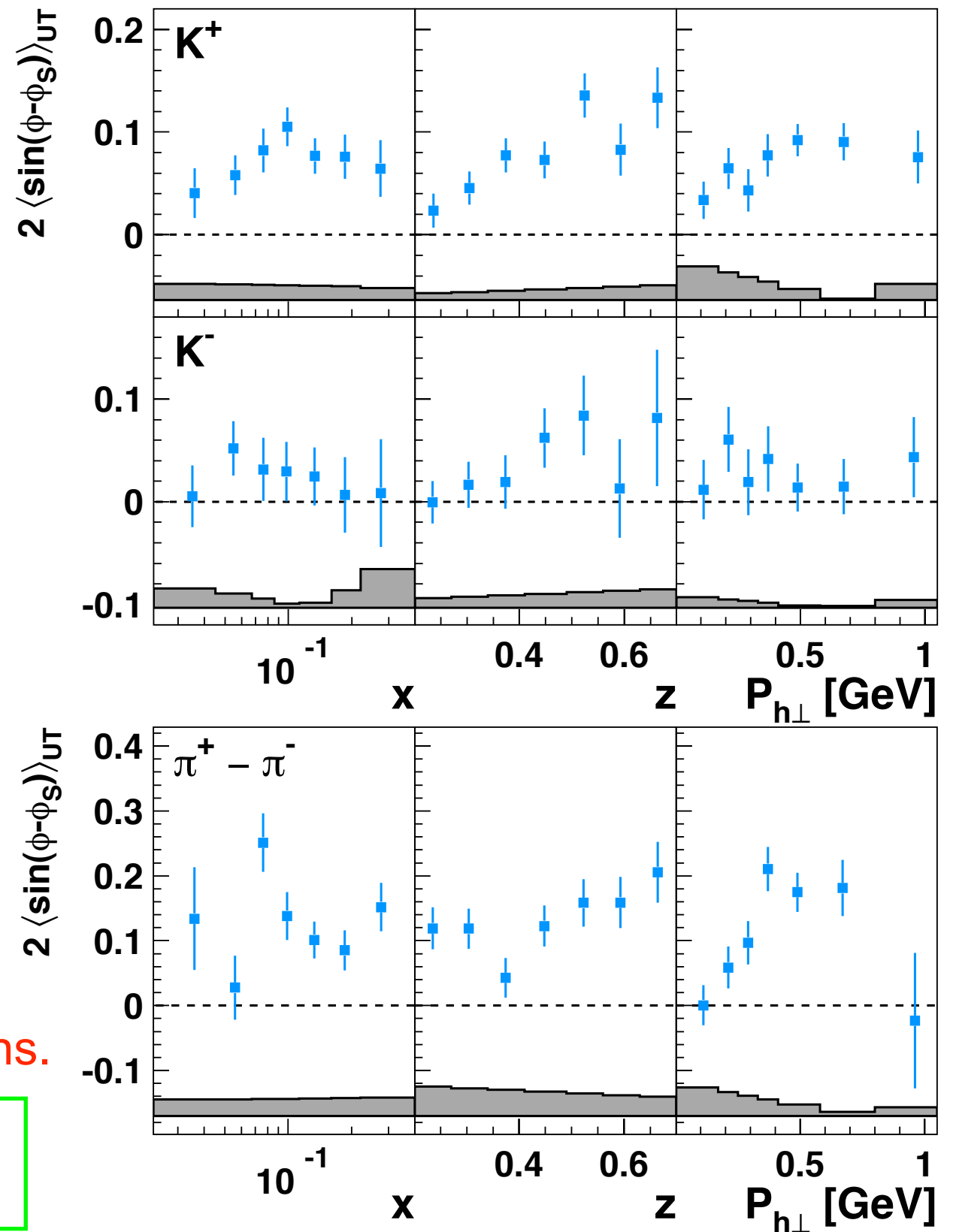
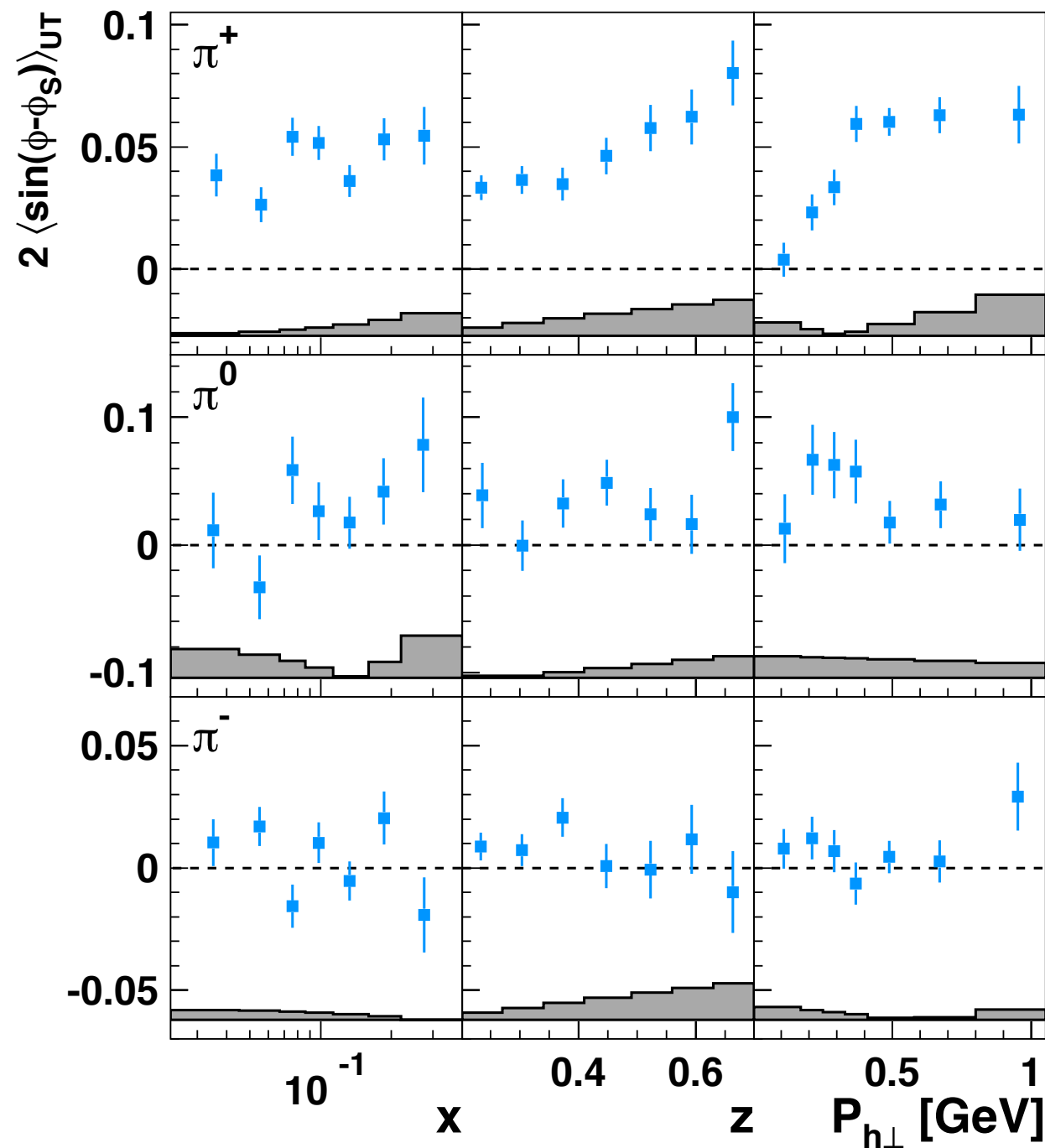
- Spin-averaged cross section *Collins, Soper, Sterman, 1985, ...*
 - Spin-dependent cross section (Sivers effect): needs further study
Boer, 2001
 - More suppressed at higher Q^2
-
- This effect (Sudakov suppression) is not included in the current formalism when extracting the Sivers functions

Collins is consistent between COMPASS and HERMES



Final published HERMES Proton: (NOT zero)

HERMES, PRL 103, 152002 (2009)

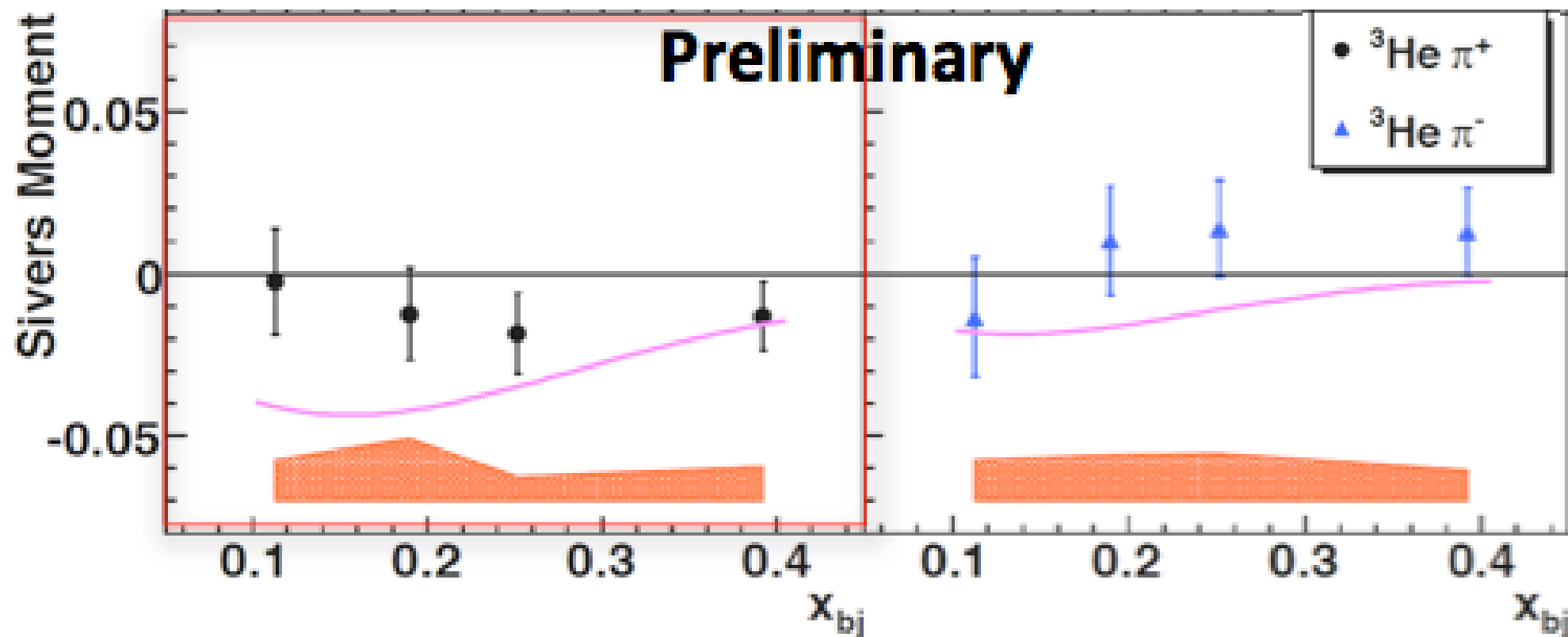


Slight change in data, main message remains.

Problem with the parametrization
of global analysis?

Sign difference between u and d Sivers functions

- u and d Sivers functions have opposite sign
 - For proton, $\pi^+ > 0$, then for neutron, $\pi^+ < 0$
 - At the same time, π^- for neutron should be smaller than π^+
 - JLAB preliminary results



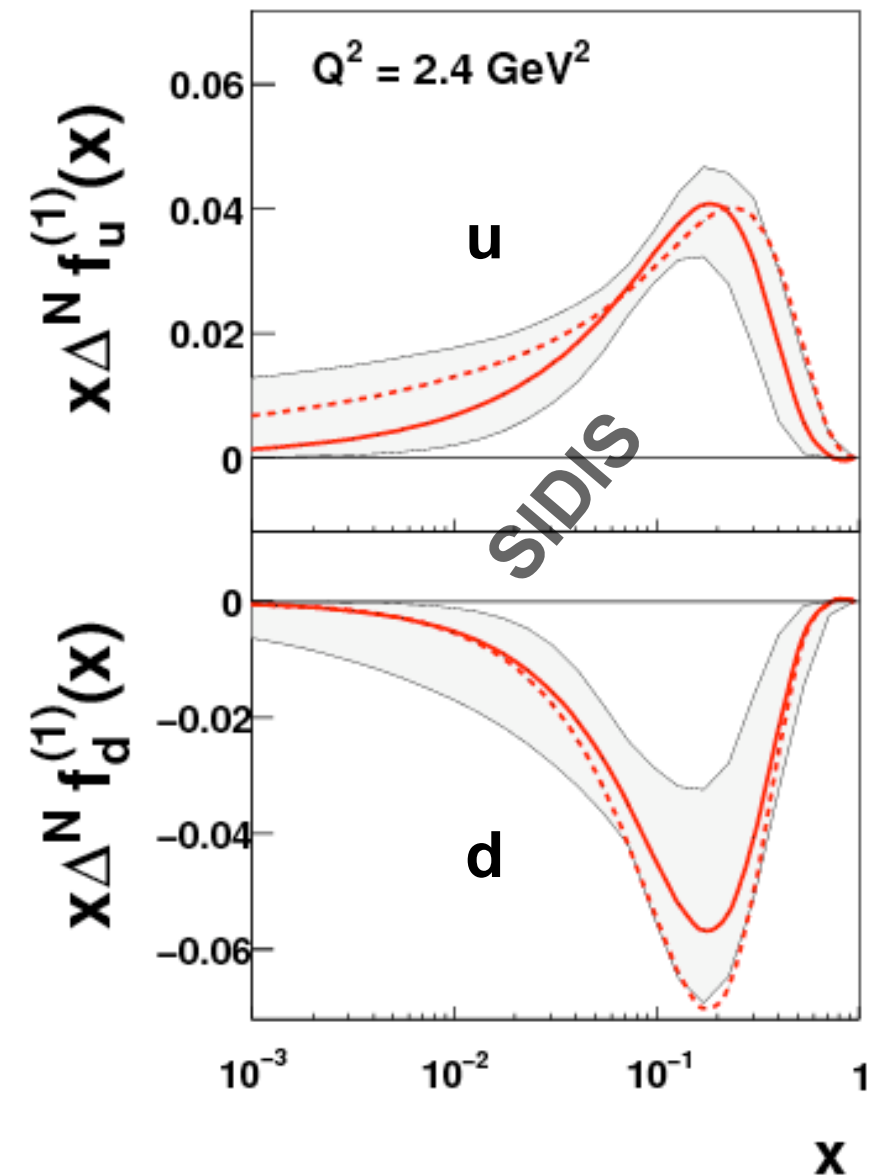
From Xiaodong Jiang's talk TMD 2010

Assume HERMES is correct

- Since theory doesn't prevent the existence of the Siver functions:

SIDIS	QCD →	DY
$\text{Sivers}_{u\text{-quark}} > 0$		$\text{Sivers}_{u\text{-quark}} < 0$
$\text{Sivers}_{d\text{-quark}} < 0$		$\text{Sivers}_{d\text{-quark}} > 0$

- u and d almost equal size, different sign
- u-Sivers is slightly smaller than d-Sivers



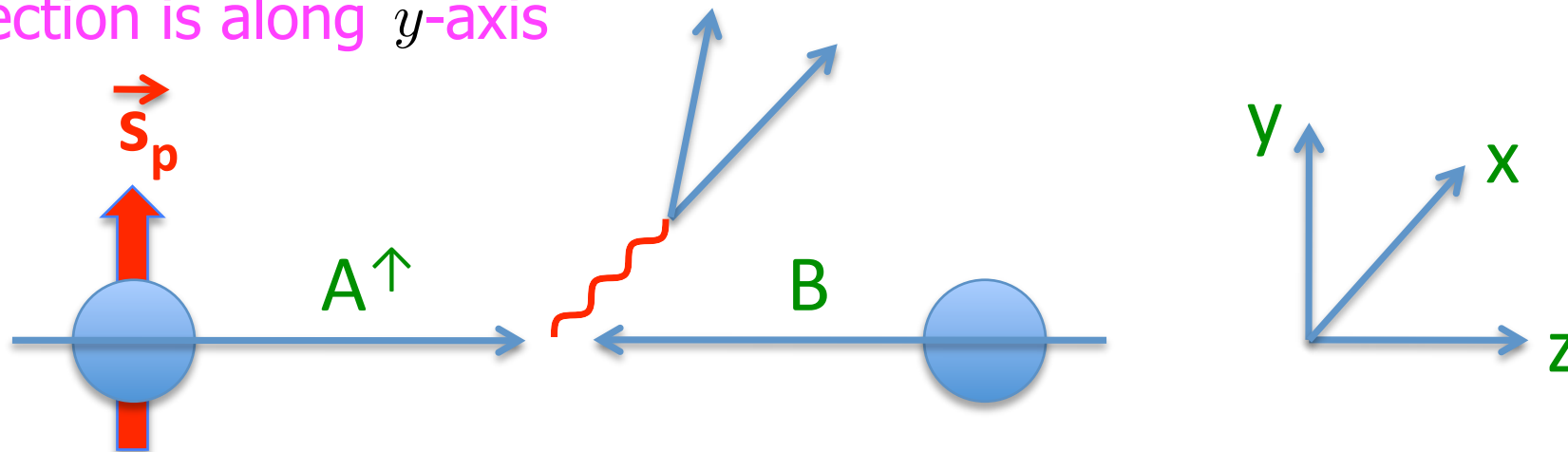
Sivers effect in Drell-Yan process

- Formula in TMD approach: weighted sum of u and d-Sivers

$$A_N = \frac{\sum_q e_q^2 \int \Delta^N f_{q/A^\uparrow}(x_1, \mathbf{k}_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})}{2 \sum_q e_q^2 \int f_{q/A}(x_1, k_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})} \propto \frac{4}{9} \Delta^N u + \frac{1}{9} \Delta^N d$$

$$\rightarrow A_N < 0$$

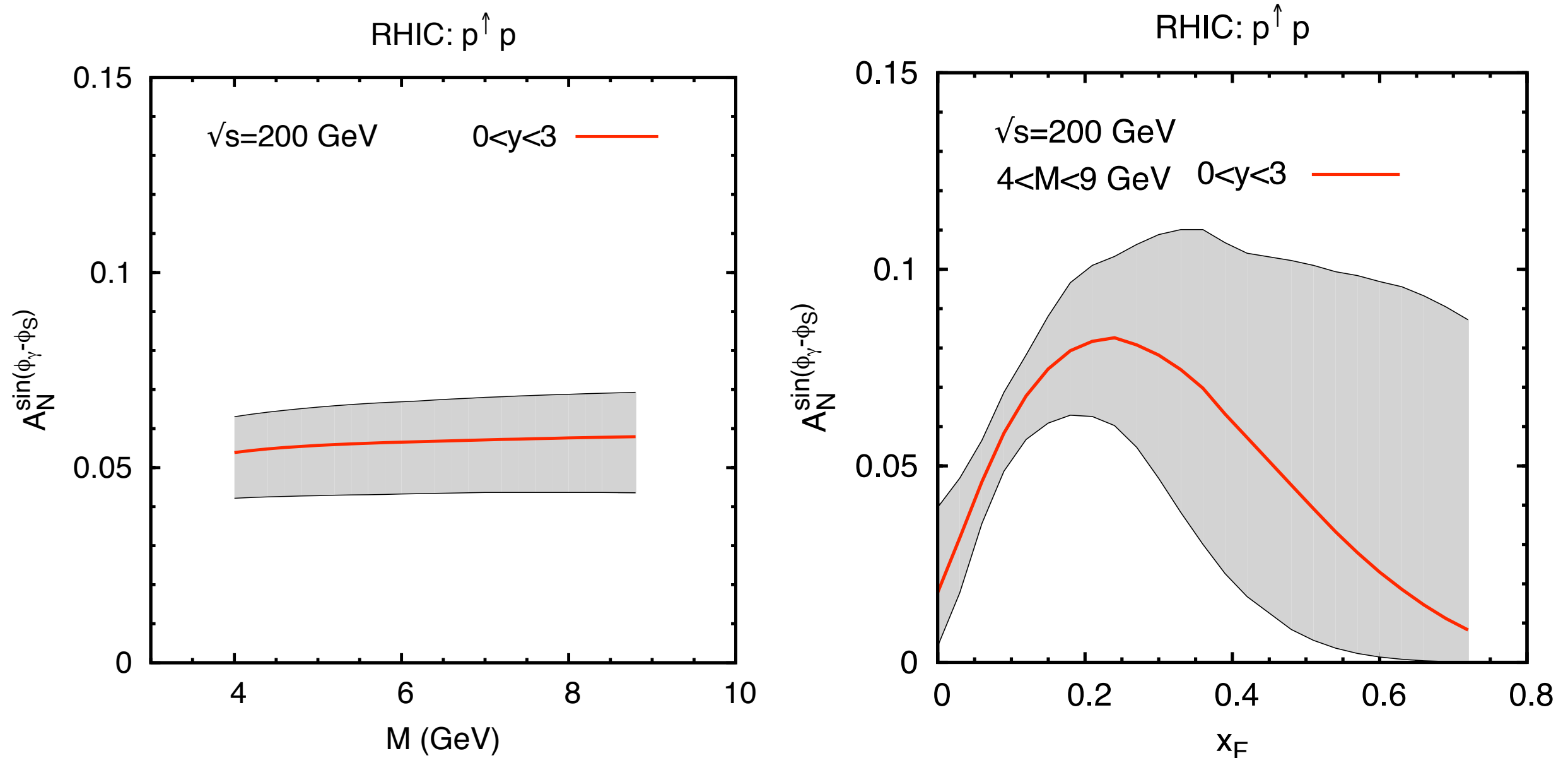
- In principle, there could also be contribution from Boer-Mulders functions times transversity: it should be very small, since it involves either anti-quark Boer-Mulders, or anti-quark transversity (or use weighted asymmetry, since they have different angle dependence)
- Careful about the frame: $A^\uparrow + B \rightarrow [\gamma^* \rightarrow \ell^+ \ell^-] + X$
 - In A-B CM frame: A^\uparrow along z -direction, B is opposite to it. “up” (\uparrow) polarization direction is along y -axis



$$\rightarrow A_N^{\sin(\phi_\gamma - \phi_s)} = -A_N > 0$$

Predictions from Anselmino's parametrizations: weighted

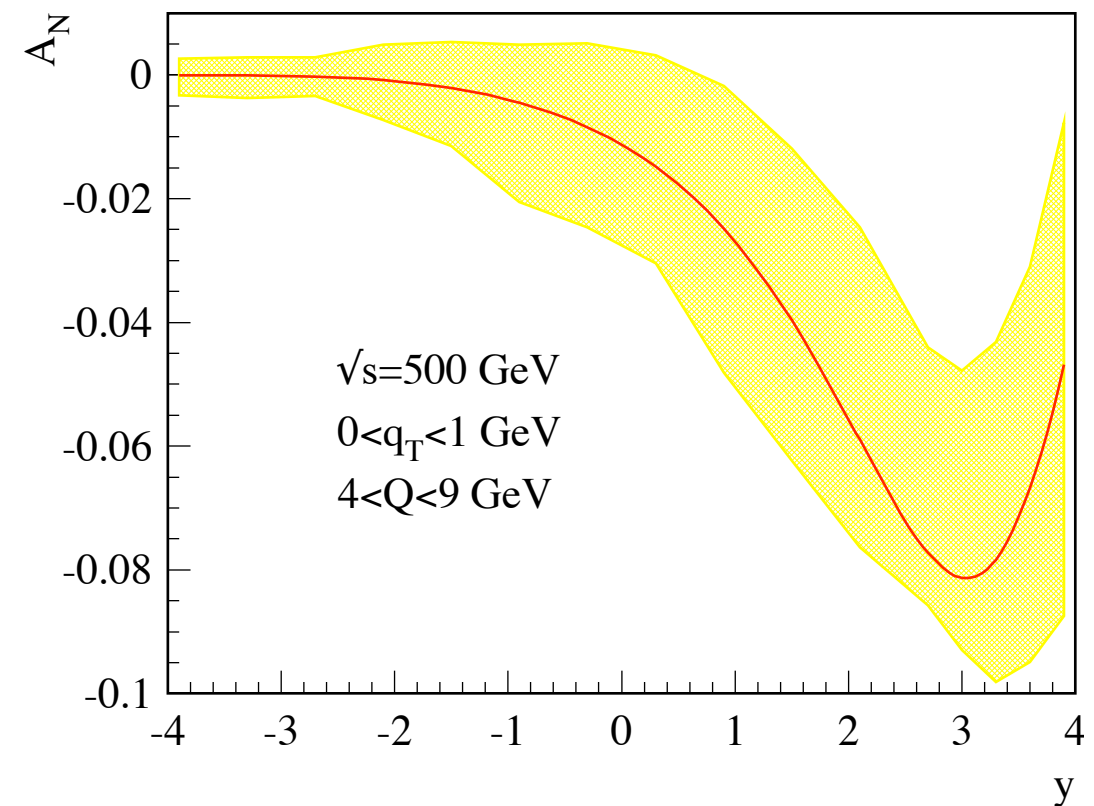
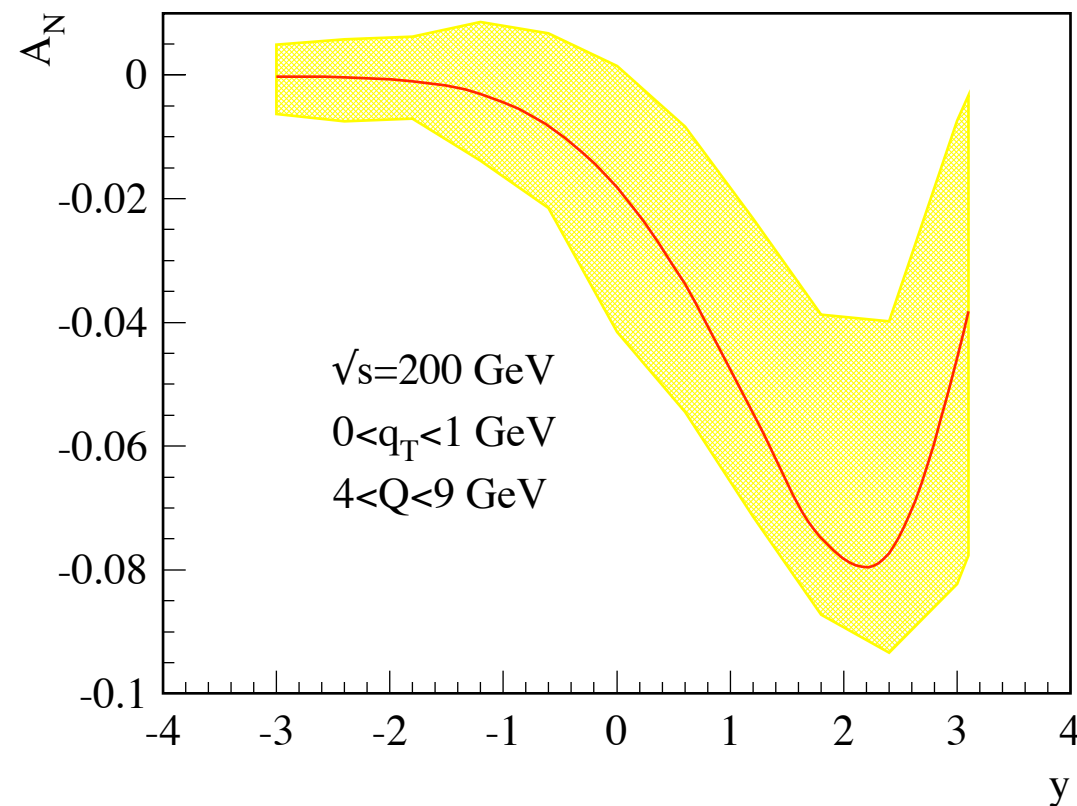
- Uncertainty band: $1-\sigma$ error of the fitted parameters in Sivers function



Anselmino, et.al, PRD79: 054010 (2009)

Rapidity dependence at 200 and 500 GeV: unweighted

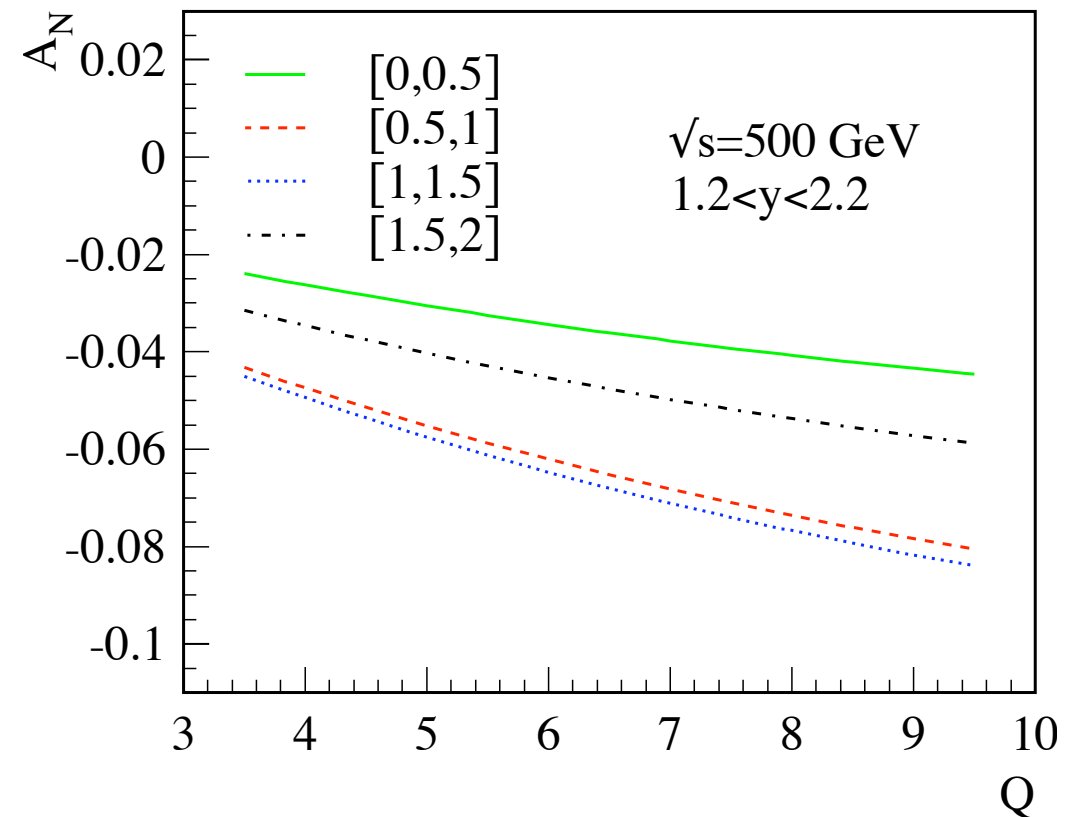
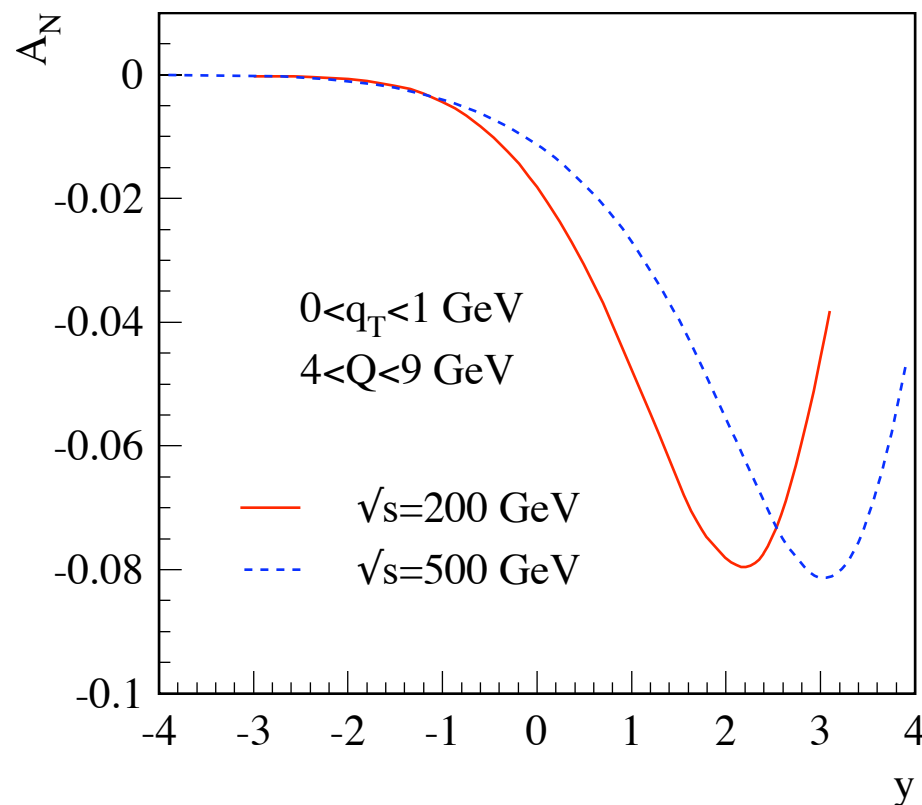
$A_N \sim 2\text{-}3\%$ in mid-rapidity $y=0$



Kang, Qiu, PRD81: 054020 (2010)

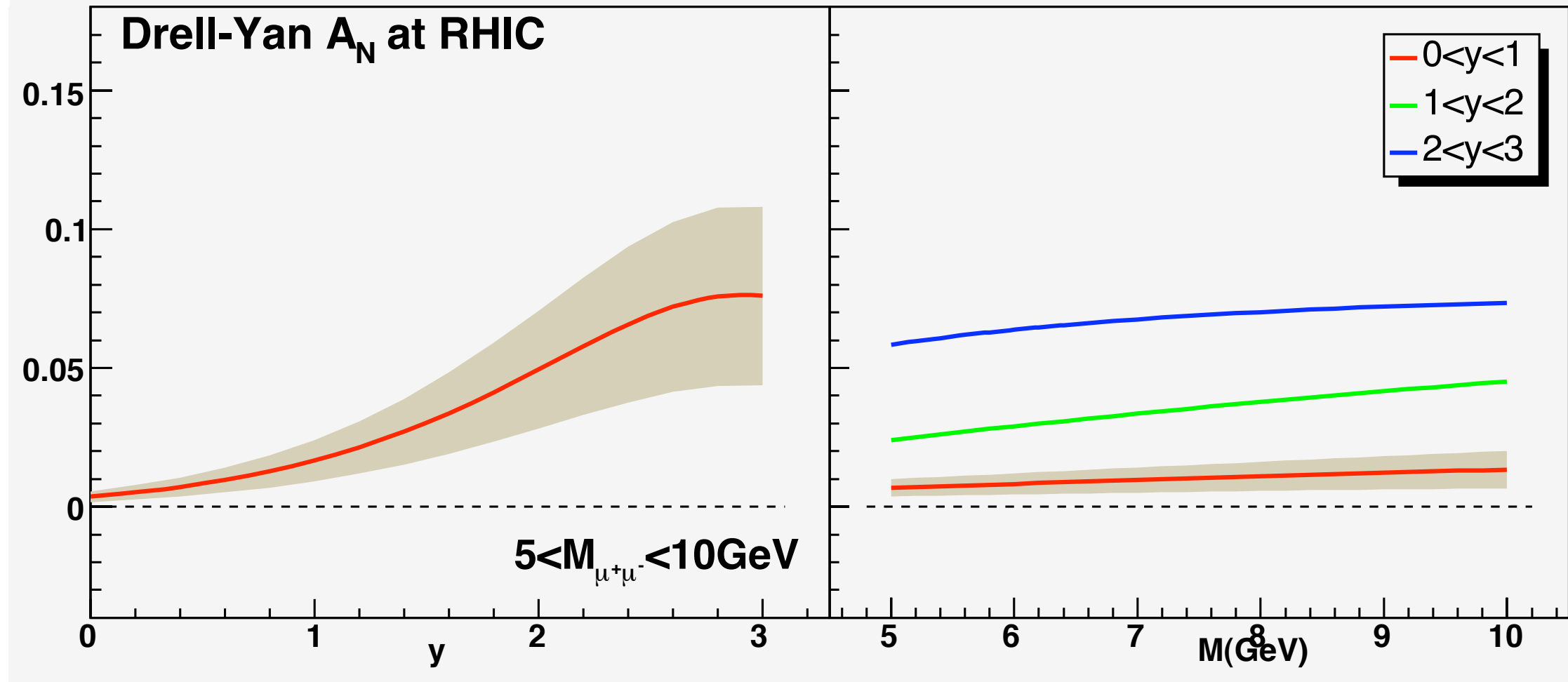
More predictions for RHIC

- Possible at RHIC at 500 GeV run?
 - 200 GeV might be difficult: not enough DY events
 - 500 GeV seems possible from simulations



Different parametrization of Sivers functions - I

- Prediction from Yuan and Vogelsang: sign convention different

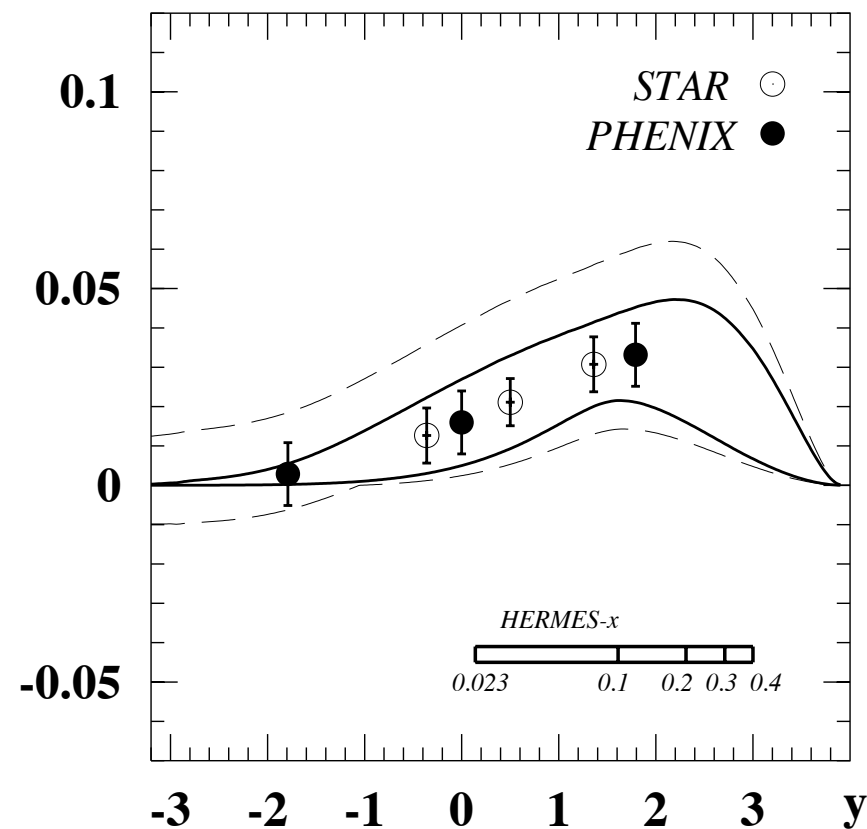


Vogelsang, Yuan, PRD72: 054028 (2005)

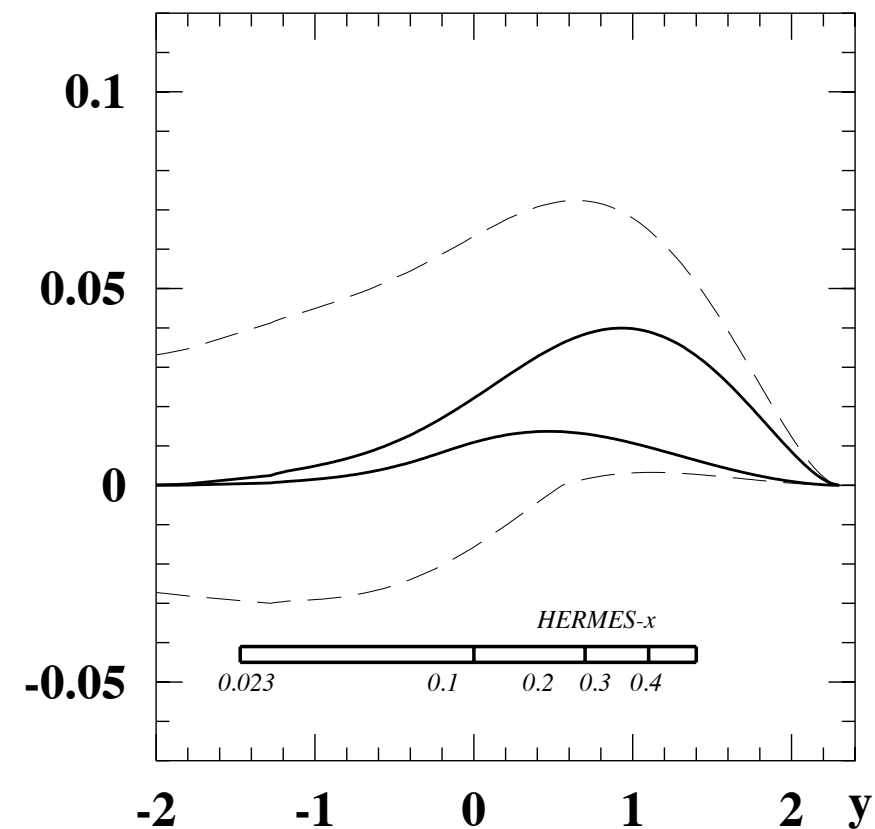
Different parametrization of Sivers functions - II

Collins, Efremov, Goeke, Menzel, et.al 2006

$A_{UT}^{\sin(\phi - \phi_S)}$ in $p^\uparrow p \rightarrow l^+ l^- X$ at RHIC $Q=4\text{GeV}$



$A_{UT}^{\sin(\phi - \phi_S)}$ in $p^\uparrow p \rightarrow l^+ l^- X$ at RHIC $Q=20\text{GeV}$



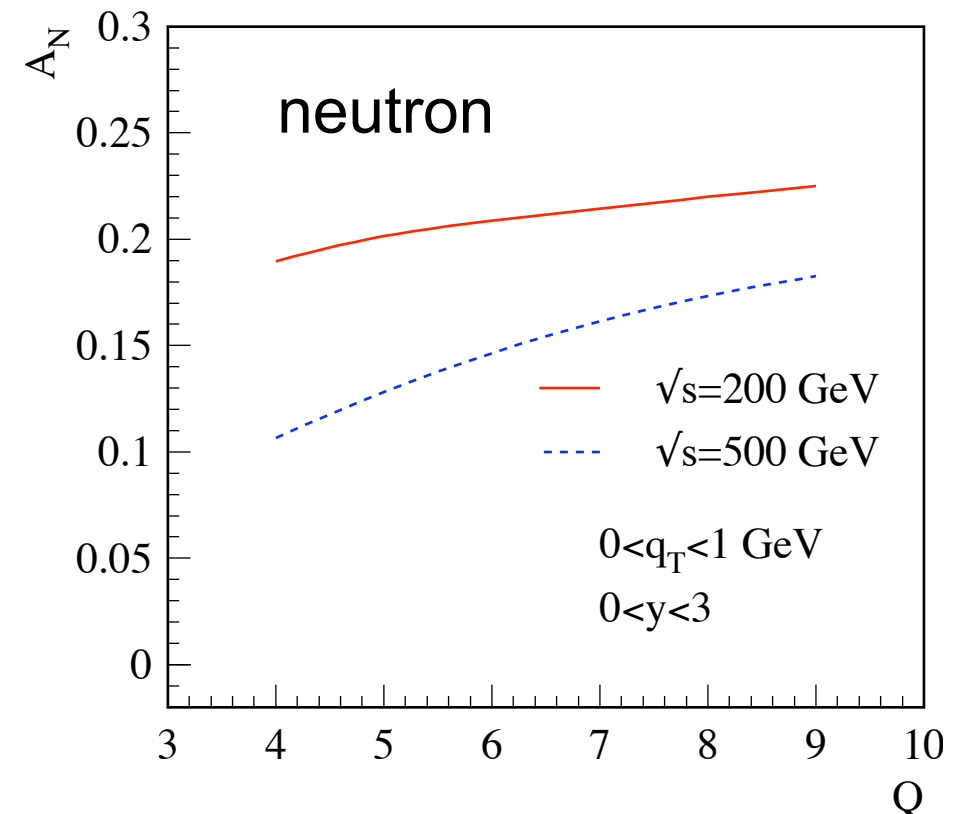
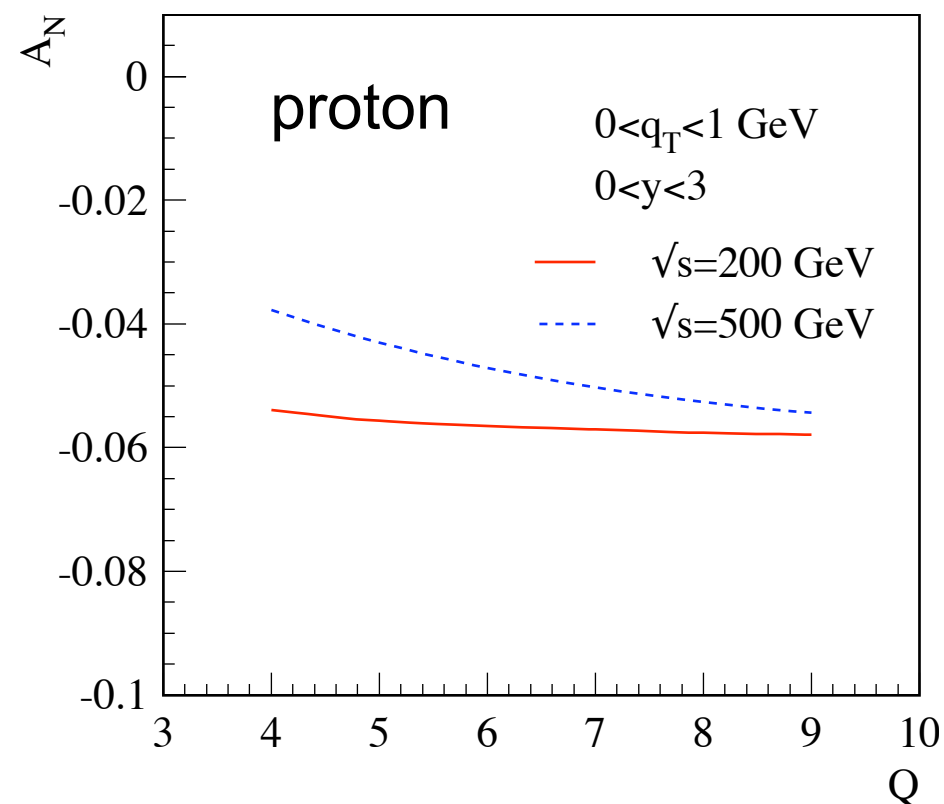
- Error band: $1-\sigma$ uncertainty of the fit of Sivers function
- Size is consistent with different parameterization

Use polarized neutron: advantage

- Sign will be opposite to the proton case

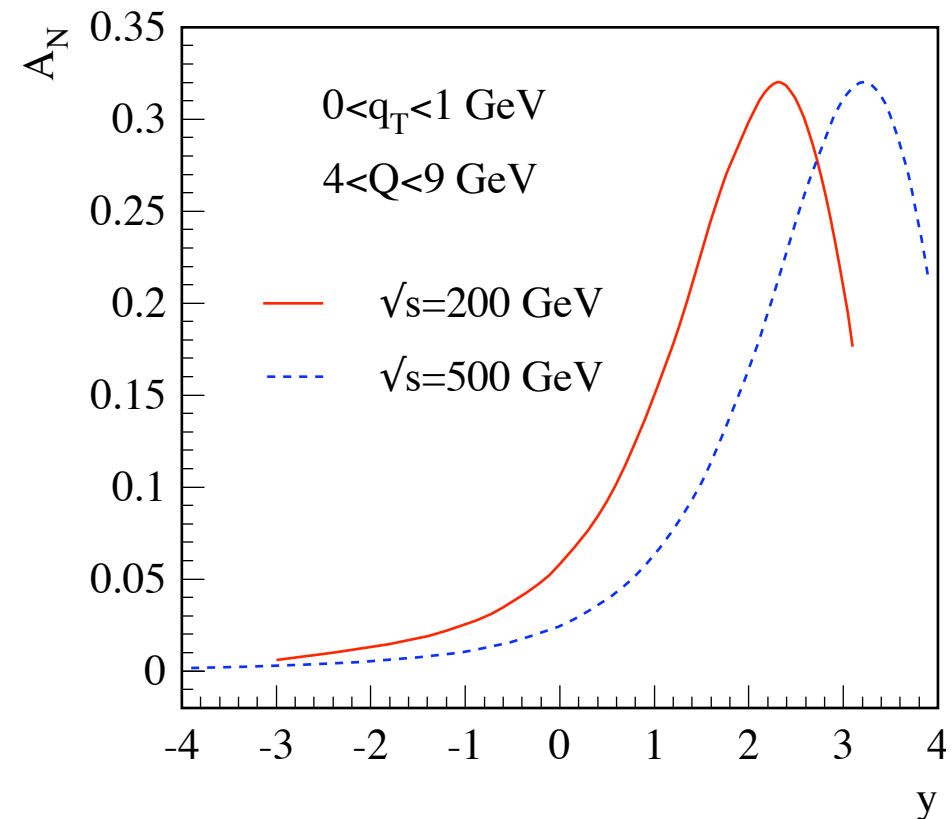
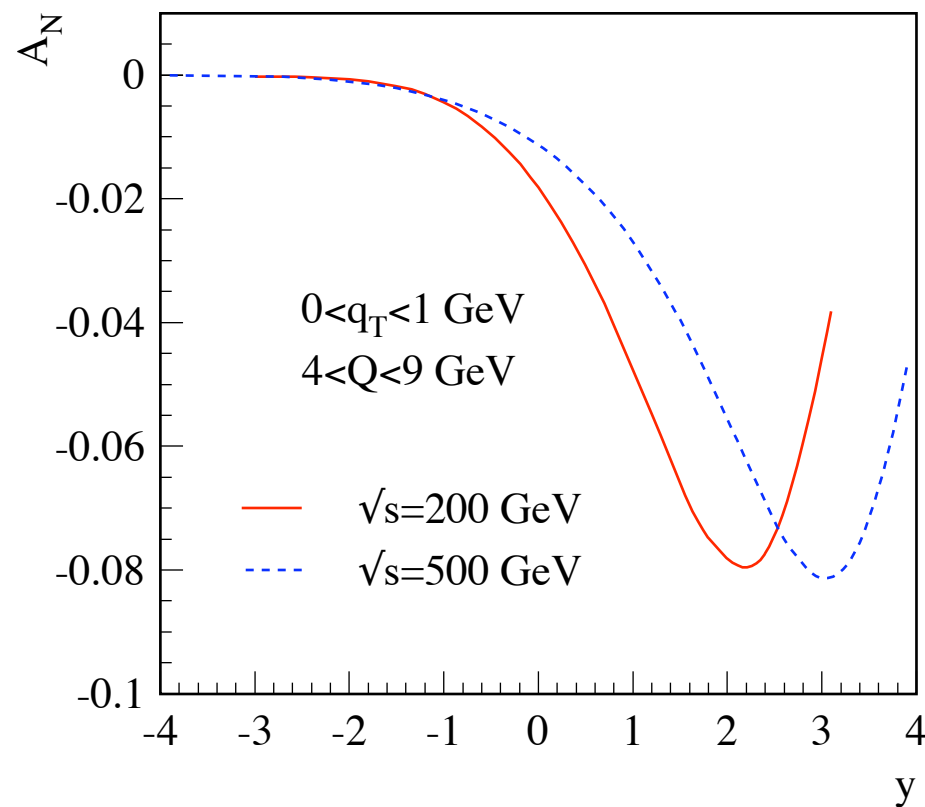
$$A_N = \frac{\sum_q e_q^2 \int \Delta^N f_{q/A^\uparrow}(x_1, \mathbf{k}_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})}{2 \sum_q e_q^2 \int f_{q/A}(x_1, k_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})} \propto \frac{4}{9} \Delta^N u_{\text{neutron}} + \frac{1}{9} \Delta^N d_{\text{neutron}} \\ = \frac{4}{9} \Delta^N d_{\text{proton}} + \frac{1}{9} \Delta^N u_{\text{proton}}$$

- d-Sivers is positive, $A_N > 0$
- d-Sivers is slightly larger, at the same time, it gets enhanced more by 4/9 compared with u-Sivers, thus the size of the asymmetry will be much bigger than the proton case



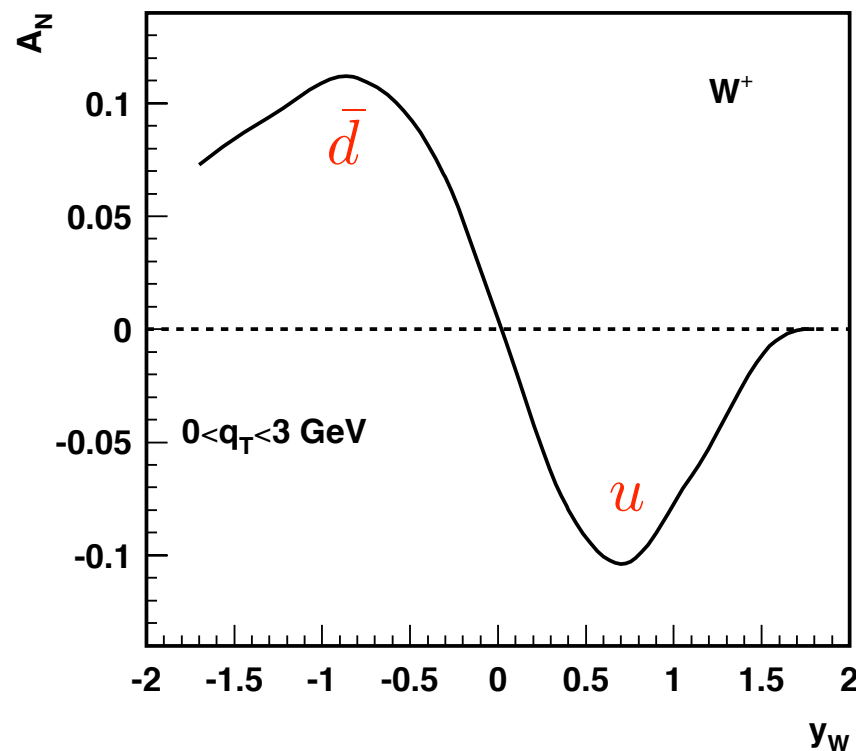
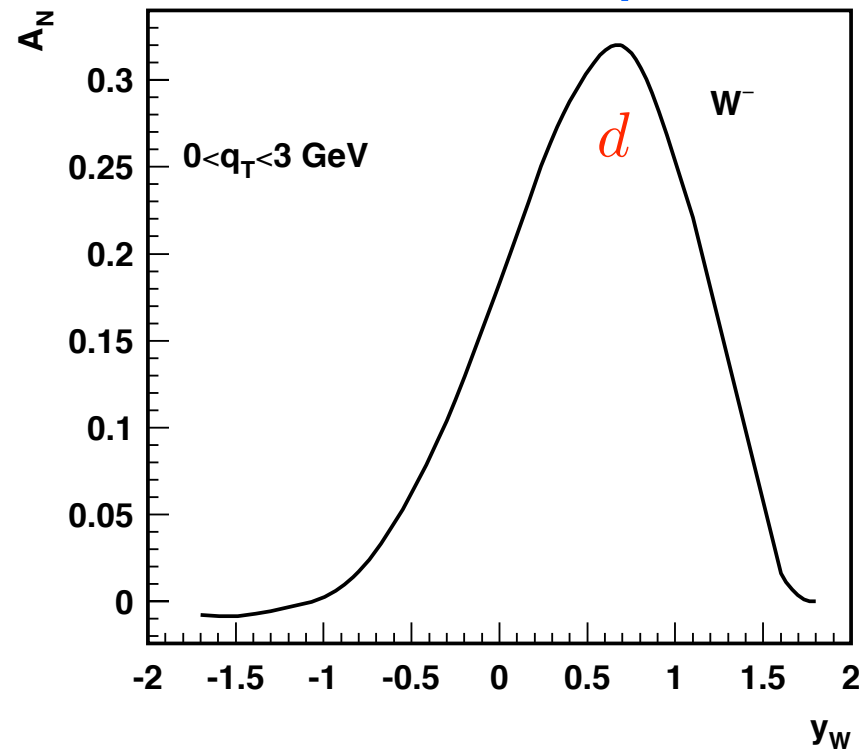
Rapidity-dependence of the asymmetry

- Positive and much larger asymmetry for He-3



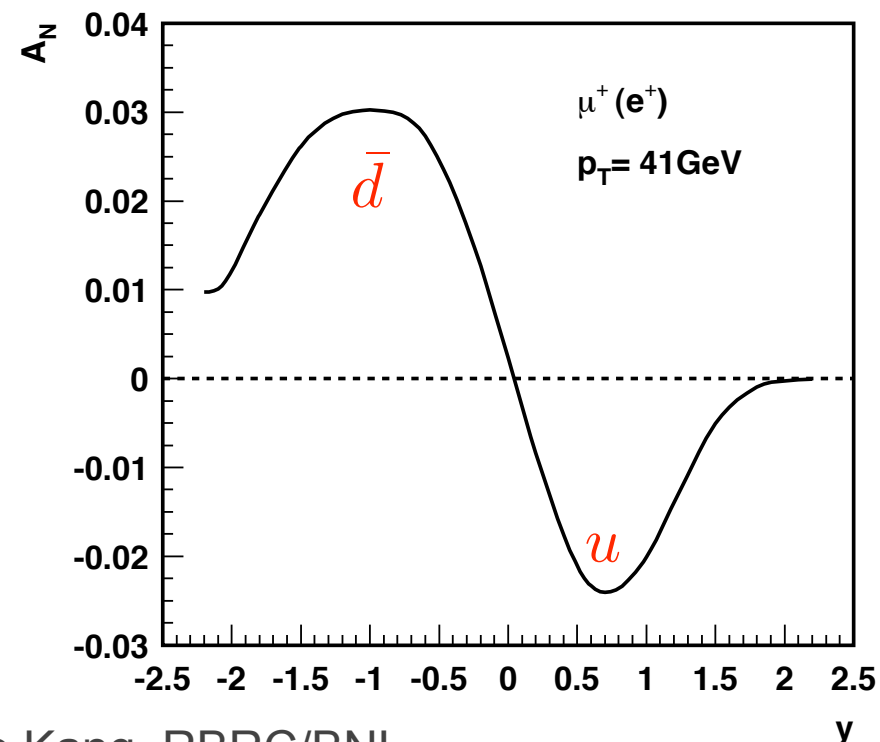
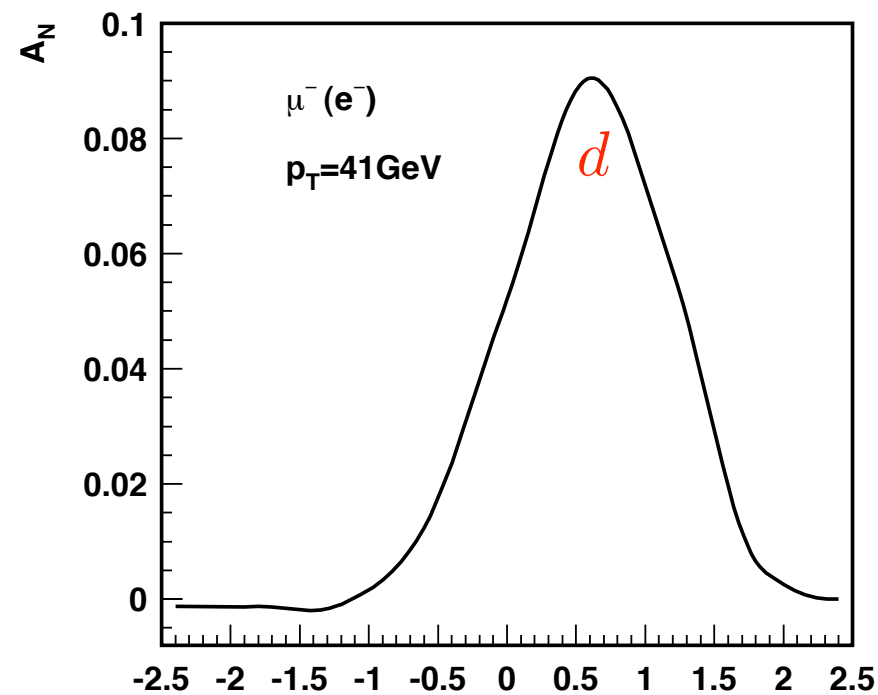
Some predictions for SSA of W bosons: pp@500 GeV

- W⁺ and W⁻ could probe different flavor of u and d Sivers function



Kang, Qiu, 2009
Zhou, Metz, 2010

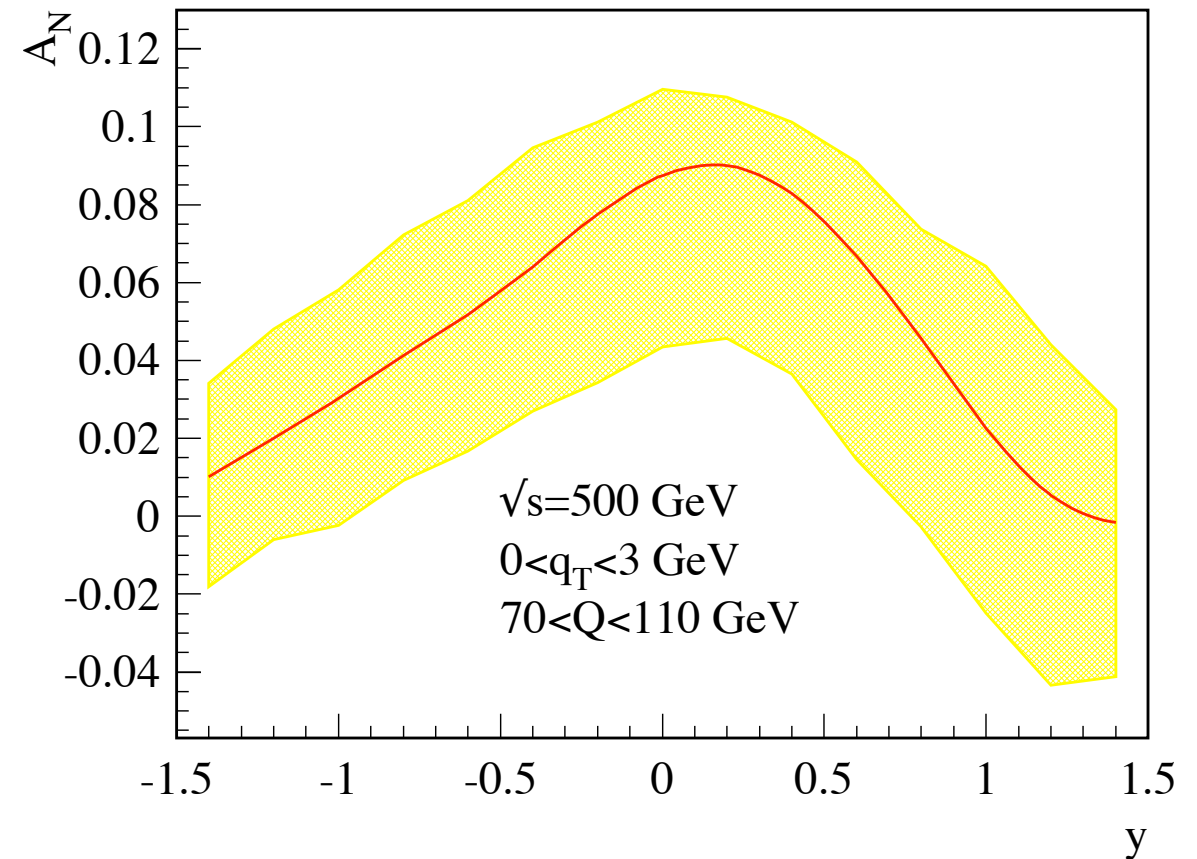
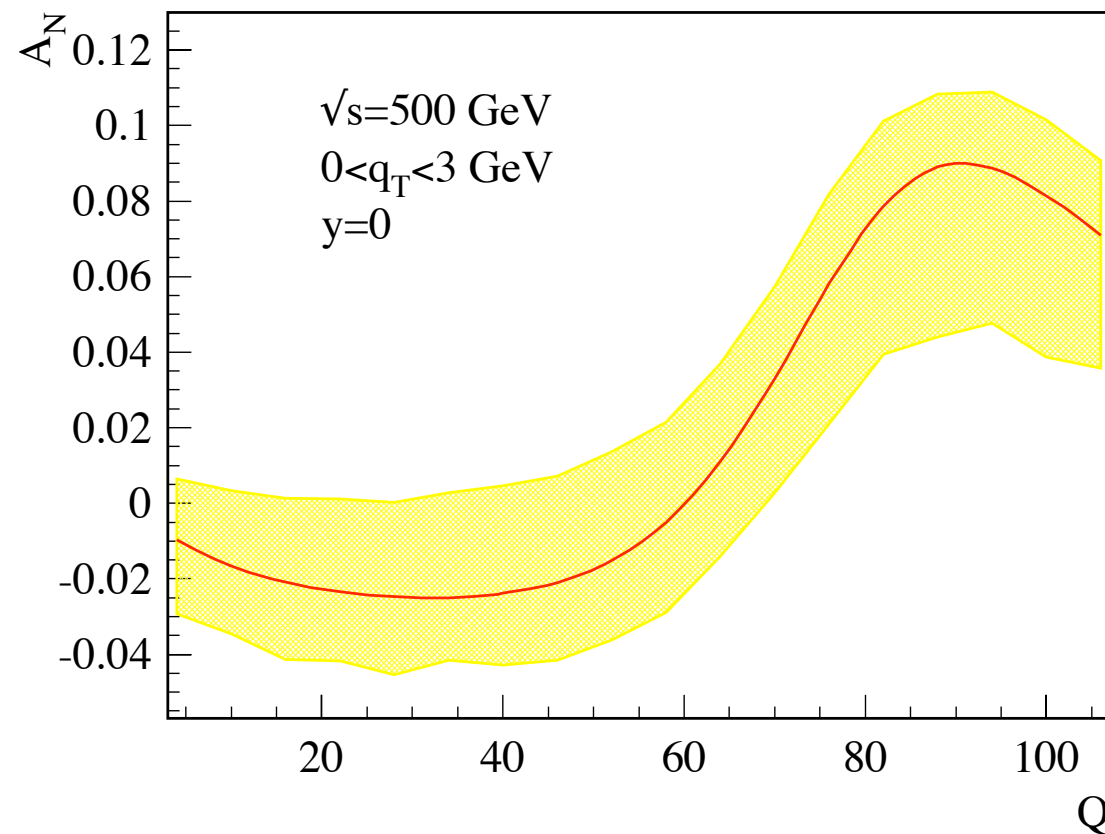
- SSA of leptons decayed from W: similar feature but diluted



SSA of Z boson at RHIC: test relative sign of u and d

- Why Z boson: change from virtual photon to Z boson, the weight of the u and d-Sivers function changes, the sign of A_N changes

Kang, Qiu, PRD81: 054020 (2010)



- Different weights: $e_q^2 \Rightarrow v_q^2 + a_q^2$

$$\begin{aligned}
 v_u &= \frac{1}{2} - \frac{4}{3} \sin^2 \theta_W & a_u &= \frac{1}{2} \Rightarrow v_u^2 + a_u^2 = 0.29 \\
 v_d &= -\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W & a_d &= -\frac{1}{2} \Rightarrow v_d^2 + a_d^2 = 0.38
 \end{aligned}$$

- The situation is not clear for W and Z productions

Other closely related measurements

- A_N for W production
 - Only measure the electron (or muon)
 - Can preserve significant sensitivity if measure just above the Jacobian peak with good p_T resolution
 - Not clear whether we can achieve sufficient resolution or not
- A_N for Z production
 - Very easy to interpret
 - **Very small cross section !!!**
 - Not clear if this is practical or not



Consequences for DY measurements

- Test the sign change of the Sivers functions between SIDIS and DY process, what are we really testing?
 - TMD approach to the SSAs
 - QCD TMD factorization
 - Our current understanding on the mechanisms of the SSAs
- If fails:
 - No sign change: our understanding for the SSAs is not complete, or not understood at all?
 - $A_N \sim 0$: Sivers is not the dominant effect for the SSAs, what's wrong with the HERMES, JLAB data? (Only COMPASS is correct?)
 - A_N is what we expect: so we are happy?
- Connection between SIDIS and pp data: talk at tomorrow afternoon

See Sivers' talk for more complete lists

- Sign change of Sivers function between DY and SIDIS is the most critical test for our current understanding of SSAs
- Let's hope we have this result as soon as possible

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Thank you